

# *Research of near-wall thermodynamic state for indoor airflow over the vertical heating unit using TIV/PIV/RTD*

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1        Research of near-wall thermodynamic state for indoor airflow  
2                    over the vertical heating unit using TIV/PIV/RTD

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8

9        **Abstract:** Up to now, few studies focus on thermodynamic state including the velocity and  
10 temperature of the air near heating unit. The thermodynamic state of the airflow over an indoor  
11 heating unit has a significant influence on indoor thermal comfort and energy consumption. This  
12 study analyzed the thermal and dynamic state of the near-wall airflow over the heating unit. The  
13 thermal state was measured using resistance thermal detectors (RTDs). The near-wall airflow field  
14 were measured by particle image velocimetry (PIV) and TIV. The performance of TIV in natural  
15 and mixed convection were evaluated by comparing the TIV and PIV measurement results. Under  
16 natural convection, the velocity shows vertical variation and the spatial difference changes more  
17 pronounced with the increase of heating temperature. Under mixed convection, the near-wall  
18 temperature changes uniform and the velocity exhibits a decreasing trend with the increase of  
19 height. Through the spectrum analysis of the temperature, it is found that the velocity measured by  
20 TIV is close to the velocity near the boundary layer to some content. The positions of the  
21 near-surface velocity measured by TIV are not fixed in all cases and change with the change of the  
22 boundary layer. The findings in this study can provide a convenient and feasible flow field  
23 measurement method suitable for actual space scale. This method can predict the effect of heating  
24 terminal units on indoor airflow and thermal environment, so as to optimize the form and  
25 arrangement of the heating terminal units, and improve heating efficiency and occupants' thermal  
26 comfort.

27        **Keywords:** Indoor velocity measurement; Near-wall airflow; Temperature distribution; Heating  
28 unit; Thermal image velocimetry

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## 29 **1. Introduction**

30       Indoors, due to natural convection, thermal plumes are formed in the vicinity of conventional  
31 vertical heating units such as panel or column radiators. In recent years, to improve the thermal  
32 efficiency of a typical vertical heating unit, some researchers imposed ventilation to heating units  
33 and analyzed the performance of the ventilation heating unit [1][2]. Under this circumstance, the  
34 flow process of the airflow over the vertical heating unit can be described as natural or mixed  
35 convection. Although the effect of influencing factors such as the surface temperature [3], the  
36 location of heating unit [4], the shape of heating unit [5], and the ventilation condition [1][2] on  
37 the distribution of indoor environmental parameters such as air temperature and air velocity in the  
38 room, have been systematically analyzed. Most of the previous studies have focused on the  
39 distributions of indoor air velocity and air temperature, and few studies have focused on the  
40 thermodynamic state near the heating unit yet. The flow and the corresponding heat transfer  
41 around the surface of heating unit are important factors that determine the distributions of indoor  
42 air velocity and air temperature. Generally, radiators are arranged against the wall under the  
43 external window. The thermal insulation of windows is insufficient compared to other envelopes,  
44 and more heat is lost through the window [6]. In addition, cold radiation caused by the window  
45 brings a certain degree of discomfort to occupants. The air adjacent to the radiator is heated and  
46 rises. When it flows through the exterior wall, heat is transferred from the air to the wall. Some  
47 heat is then lost from the exterior wall. A part of the heated air flow to the inner of the room due to  
48 the insufficient thermal buoyancy and exchange heat with the other air and surfaces of the room.  
49 Therefore, the thermodynamic state of the air near the radiator directly affects the heat transfer  
50 process among the air and surfaces in the room and the discomfort brought by the cold window

51 radiation. Precise estimation of the thermodynamic state of the airflow over the heating unit in a  
52 room is of great significance for better understanding the thermal characteristics of an indoor  
53 heating system [7][8][9].

54 To evaluate the thermodynamic state of the airflow over a heating unit, accurate measurement  
55 is necessary. The current air temperature measurements are mainly conducted using the resistance  
56 of thermal detectors (RTDs) or thermocouples. The air temperature measurement can be easily  
57 performed owing to the high frequency, short response time and the ability to store large amount  
58 of data of the temperature sampling instruments. On the other side, the velocity measurement  
59 methods are mainly divided into two kinds: the point-wise and the globe-wise measurement  
60 methods [10]. The point-wise measurement method can only reflect the velocity of the  
61 representable points instead of the spatial distribution, and the measured data of measurement  
62 points is mostly used as validation or boundary condition of numerical simulation. The globe-wise  
63 measurement method can visualize the flow field [11]. Particle image velocimetry (PIV) is a  
64 typical and the most widely used flow visualization method. The adjustment of laser and, camera,  
65 tracer particle concentration, and image processing all require certain expert knowledge [12].  
66 Moreover, PIV measurement system components are complex and fragile, and thus not suitable for  
67 movement; therefore, this method is mostly used in laboratory research. Except some special cases  
68 [13][14], traditional velocity measurement methods are rarely applied successfully in the flow  
69 field measurement near the indoor heated surface. Besides the complexity of the indoor  
70 environment, one of the main reasons for this is the limitation of various velocity measurement  
71 technologies, which are explained in detail in [15].

72 The contradiction between the limitations of existing point-wise or globe-wise measurement

73 methods and the need to measure specific surface velocity has been increasing, and researchers are  
74 now seeking other velocity measurement methods. The interaction of surface turbulent structures  
75 with surface temperature fluctuations has been draw attention gradually. The turbulent structures  
76 sweep the interface and give rise to surface temperature fluctuations [16][17]. Recent studies have  
77 revealed that the characteristic of the near-surface airflow can be indirectly deduced based on  
78 surface high frequency temperature fluctuations. In terms of the acquisition of high frequency  
79 temperature fluctuations, thermal infrared camera is more advantageous owing to its high  
80 sensitivity, short response time, and spatial resolved capability, in comparison with thermocouple  
81 or resistance temperature detectors [18]. As a result, infrared thermography is gradually used as an  
82 alternative tool for evaluating near-surface flow distributions. In the current study, the use of  
83 infrared thermography in obtaining the characteristic of the near-surface airflow is mainly based  
84 on three methods. a) Theoretically establishing the correlation between near-surface velocity and  
85 temperature fluctuations based on the surface renewal (SR) theory [16]. The near-surface flow  
86 structures are composed of eddies with various sizes that fluctuate with various frequencies and  
87 continuously eject and sweep back to the surface. As continuously interacting with the interface,  
88 the near-wall flow structures leave traces in the form of surface temperature fluctuations. In  
89 addition, the near-surface eddies are distributed according to specific distribution laws [19]. By  
90 exacting the eddy distribution parameters from the high frequency surface temperature  
91 fluctuations, in combination with the relation between the surface velocity and near-wall heat  
92 transfer, the near surface velocity can be deduced [19]. b) The other one is noted as TIV (thermal  
93 image velocimetry). This method is similar to the PIV method, which estimates the near-surface  
94 velocity by analyzing the displacements of the surface thermal spots within consecutive frames

95 captured by the infrared thermal camera. The basic principal and preliminary application of this  
 96 method can be referred in [15]. c) The last method is also based on the surface renewal theory.  
 97 Unlike the first method, the near-surface turbulent state is deduced according to the correlation  
 98 between the surface temperature fluctuations and the intermittent near-surface eddies and no  
 99 quantitative velocity is acquired.

100 **Table. 1.** Application of infrared thermography in near-surface velocity measurement

Object	Indoor/Outdoor	Scale model/ Field	Methods	Velocity Range	Reference
Horizontal plate	Indoor	Scale model	(a)	0.6 - 4 m/s	[19]
Evaporation plate	Indoor	Scale model	(a)	0 - 0.5m/s	[20]
Vertical plate	Indoor	Scale model	(b)	0 - 0.3 m/s	[15]
River	Outdoor	Field	(b)	0.9 - 1.1m/s	[21]
Volcano	Outdoor	Field	(b)	—	[22]
Solar tower	Outdoor	Field	(b)	—	[23]
River	Outdoor	Field	(b)	0 - 4m/s	[24]
Wall	Outdoor	Field	(b)	0 - 2m/s	[25]
Horizontal plate	Outdoor	Scale model	(b)	0 - 0.005m/s	[26]
Lava	Outdoor	Field	(b)	0 - 0.006m/s	[27]
Artificial turf ground	Outdoor	Field	(c)	—	[28]
Vineyard	Outdoor	Field	(c)	—	[29]
Ground	Outdoor	Field	(c)	—	[30]

101 1. (a), (b) and (c) denote the three methods described above.

102 2. Scale model and field denote the scale model experiment and the field measurement

103 Table. 1 summarizes the current main applications of infrared thermography in near-surface  
 104 velocity measurement. The first method obtains the near-surface spatial-averaged velocity. For  
 105 some specific situations, especially the flow field above the heat source, the spatial-averaged flow  
 106 velocity cannot reflect the spatial distribution characteristics of velocity, and thus the application  
 107 of this method is limited. The third method can only provide the qualitative velocity and is not  
 108 suitable for the evaluation of velocity distribution. Compared with the other two methods, the  
 109 second method can acquire the near-wall two-dimensional velocity distribution and is the most  
 110 widely used. As shown in Table. 1, TIV is mostly used in outdoor research at present. For both the

111 indoor and outdoor environment, the primary condition for applying TIV is that the target surface  
112 should guarantee certain temperature fluctuations. In outdoor studies, especially for targets such as  
113 lava or solar tower, the temperature fluctuation is intense, which is conducive to the application of  
114 TIV. Nevertheless, there are still some conditions that cannot meet the temperature fluctuation  
115 requirement. In some outdoor studies [25], a blackened polystyrene board was pasted on the  
116 observed surface when the target surface temperature fluctuation was small. The heat capacity of  
117 the polystyrene board is small and the blackened surface absorbs more solar energy. Through  
118 using this method, the surface temperature fluctuations increase greatly. Compared with the  
119 outdoor applications, the utilization of TIV in indoor research is more difficult because of the  
120 small temperature fluctuations that occur in indoor environments. In the micro-environment over  
121 the heating unit, the air near the heating unit is heated and rises vertically under the buoyancy  
122 force. In addition, the vertical wall near the heating unit is generally colder than the heated air, and  
123 heat is transferred from air to the wall [3]. Owing to the combined effect of the vertical cold wall  
124 and the heating unit, the unique thermodynamic state of the airflow is formed. If the black-painted  
125 benzene board is placed on the cold wall, it will inevitably affect the original heat transfer process,  
126 and make the result deviate from the actual flow to some extent. In addition, most velocity  
127 information obtained in the references summarized in Table. 1 is not validated by the other kinds  
128 of mature measurement methods, and the precision of the results measured by TIV cannot be  
129 guaranteed. Therefore, TIV is in its early stages of development, and its considerable potential is  
130 unexplored.

131 In our earlier work, we proposed a TIV-based method for indoor environments and verified  
132 the feasibility of the method for velocity measurement under natural convection [15]. However,



133 the thermodynamic state of the airflow over the heating unit and the relationship between the  
134 thermodynamic state of the airflow and heating temperature and ventilation condition has not been  
135 extensively explored yet. This is a subsequent study of [15] and the following questions are  
136 explored:

- 137 1. The performance of TIV under natural convection and mixed convection.
- 138 2. The characteristic of the thermodynamic state of the airflow near the heating unit when the  
139 heating temperature changes or mechanical ventilation is imposed.

## 140 **2. Methodology**

### 141 **2.1. Experimental setup and procedure**

142 Experiments were performed in a closed rectangular cavity shown in Fig. 1. The coordinate  
143 axes are shown in Fig. 1. The cavity was placed in an air-conditioned room whose temperature  
144 was maintained at 20 - 21 °C. The cavity had an internal dimension of 1200 (L) × 500 (W) × 1400  
145 (H) mm. The walls were made of transparent Plexiglas. The emissivity of the cavity surfaces was  
146 estimated using the infrared thermal camera and a calibrated thermocouple, and was equal to 0.90.  
147 An electric heating unit of which the dimension was 500 (W) × 400 (H) mm, was attached to the  
148 lower part of the vertical wall as a vertical heating element. A 5 - mm - thick polyvinyl chloride  
149 with a thermal conductivity of 0.2 W/ (m·K), was stuck between the heating unit and the vertical  
150 wall to decrease heat conduction from the heating unit to the back of the vertical wall. The cavity  
151 was connected to the air duct, and a fan provided mechanical ventilation from the ventilation inlet  
152 positioned at the bottom of the cavity. The size of the ventilation inlet and outlet was 500 (W) × 20  
153 (H) mm. The heating unit was under the control of an electronic thermostat. The experimental

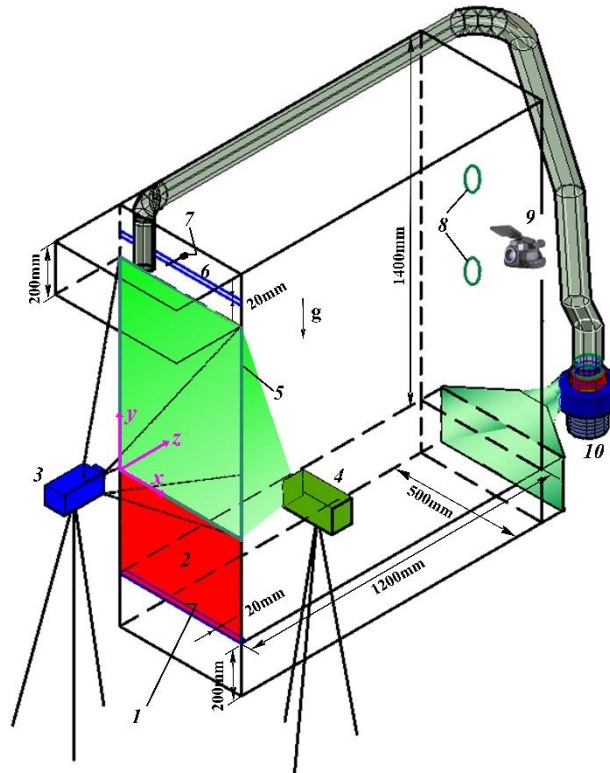
154 cases, shown in Table. 2, were set by adjusting the heating temperature of the heating unit and the  
 155 ventilation velocity.

156 Before the temperature and velocity measurements can be taken, it takes six hours in advance  
 157 to reach a steady state for temperature and flow fields. As shown in Fig. 1, one RTD mounted at  
 158 the ventilation outlet was used to measure the air temperature continuously. Once the sampled  
 159 temperature stabilized and the temperature fluctuation was less than 0.1 °C /min [31], it was  
 160 assumed that a stable thermodynamic state within the cavity has been reached, and the  
 161 measurements were taken.

162 **Table. 2.** Experimental cases

	Surface temperature of heating unit (°C)	Ventilation velocity (m/s)	Heat transfer form
Case 1	45	0	Natural convection
Case 2	65	0	Natural convection
Case 3	65	1.75	Mixed convection

163



164

165

**Fig. 1.** Experimental setup: 1. Inlet 2. Heating unit 3. CCD camera 4. Laser 5. Visualized surface

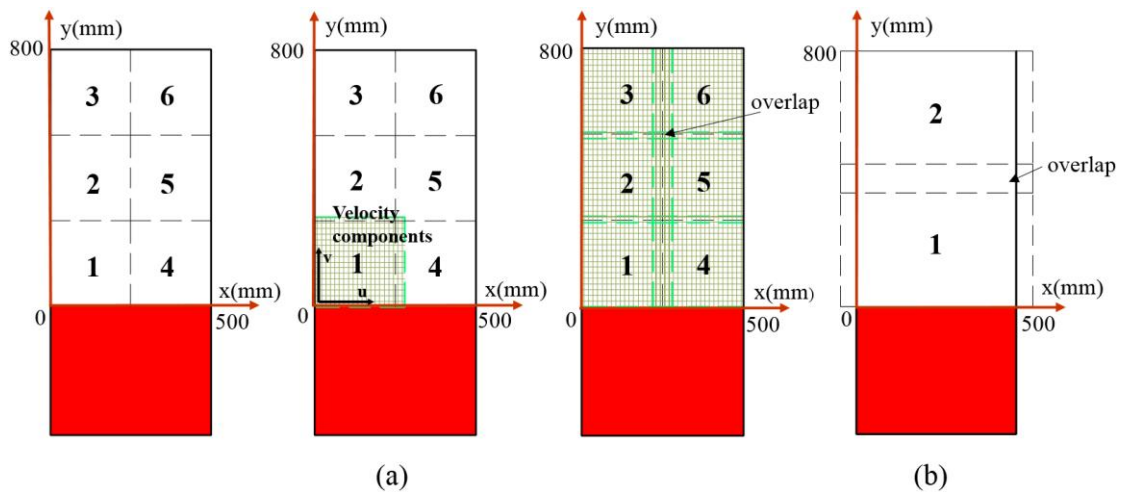
166

6. Outlet 7. RTD 8. Shooting positions 9. Infrared camera 10. Fan

167 **2.2 Measurements and data processing**

168 2.2.1 PIV measurement

169 As shown in Fig. 1, PIV was used to visualize the two-dimensional airflow fields that are  
 170 8mm and 15mm away from the vertical surface behind the heating unit. The entire visualized area,  
 171 which is the region above the heating unit shown as Fig. 1, is 500 (W) × 800 (H) mm. As shown as  
 172 Fig. 2 (a), the dimension of the view field is 280 (W) × 280 mm (H) and there are overlap areas  
 173 conducive to connecting the separate measurement domains. The six airflow fields are connected  
 174 using the technique stated in [32]. The relevant parameters of PIV system are shown in Table. 3.



175

176

177

**Fig. 2.** Schematic of measurement domains a. PIV b. TIV

**Table. 3.** PIV parameters

Name	Parameters
Laser model and power	Double-pulsed Nd:YAG laser, 150mJ/pulse
CCD model	PIVCAM13-8
Lens model	AF Nikko 50mm f/1.8D Lens
Visual field of CCD	2048pixel × 2048pixel
Size of the interrogation window	32pixel × 32pixel, 50% overlap
Dimension of the view field	280 (W) × 280 (H) mm
Spatial resolution	0.137mm/pixel
Sampling frequency	7Hz
Sampling pairs	600

## 178 2.2.2 TIV measurement

179 The accuracy of the measurement results cannot be guaranteed when the infrared camera  
180 captures the temperature fluctuations of the target surface through the Plexi-glass. In addition, the  
181 tracer particles suspended in space generate a considerable amount of scattering interference.  
182 Considering these two factors, the CCD camera and the infrared camera were positioned at the  
183 opposite sides. In this way, the measurement results of TIV are mirrored to make comparison with  
184 the results of PIV. In each shooting position, a circle hole fit for the camera lens was dug. The two  
185 holes were covered with lids to prevent the inner air from escaping. When conducting the TIV  
186 measurement, the lid was removed and the infrared camera captures the surface temperature of the  
187 observed Plexi-glass surface without the obstacles of any surfaces. The tiny space between the  
188 camera lens and the circle hole is sealed with cotton felt to stop the air leakage. Because the  
189 camera was placed and removed in a very short time and the shooting hole is very small, the  
190 amount of the gas entering or exiting the cavity at the moment of removal or resetting the lid can  
191 be negligible.

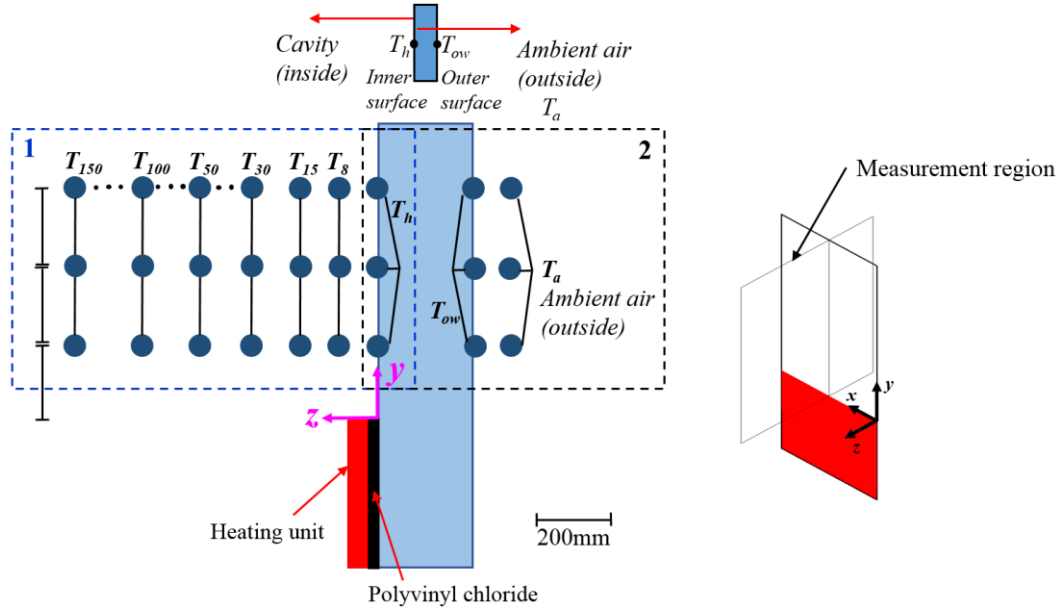
192 The surface temperature of the vertical wall was captured by a FLIR-T1040 infrared camera.  
193 Two *.ats* video files were acquired in two shooting positions located at 600 mm and 950 mm away  
194 from the bottom of the cavity. The entire visualized area is 500 (W)  $\times$  800 (H) mm and  
195 decomposed into two domains and each domain area is 598 (W)  $\times$  445 (H) mm, shown as Fig.2  
196 (b). The two airflow fields were connected with the method used in PIV. The detailed information  
197 of the parameters for thermal image velocimetry is shown as Table. 4.

198 **Table. 4.** Infrared camera parameters

Name	Parameters
Model	FLIR T1040
Resolution	1024pixel × 768pixel
Sensitive wave length range	7.5~14μm
noise-equivalent temperature difference	25mK
Angel of version	28° × 21°
Dimension of the view field	598 (W) × 445 (H) mm
Spatial resolution	0.584mm/pixel
Sampling frequency	30Hz
Sampling duration	20s

199 2.2.3 Temperature measurement

200 The temperature measurement was carried out using 32 RTDs. As placing so many RTDs in  
201 the cavity may disturb airflow field, the temperature measurements and velocity measurements  
202 were carried out separately. The RTDs were calibrated by using a set of temperature calibration  
203 system. The RTDs were calibrated in the temperature range of 18 to 46 °C, at fifteen points with  
204 an interval of 2 °C. The linear calibration coefficients of the RTDs were acquired by regression  
205 analysis. The maximum error of the 33 RTDs was 0.347 °C before calibration. According to the  
206 linear correction of each RTD, the maximum error was decreased to 0.123 °C. The temperature  
207 measurement consists of two aspects. One is to investigate the temperature distributions inside the  
208 cavity, shown as the dotted square labeled “1” in Fig. 3. The other aspect is to evaluate the inner  
209 heat transfer of the vertical wall and the measurement points are shown in the dotted square  
210 labeled “2” in Fig. 3.



211

212

**Fig. 3.** Schematic distributions of temperature measurement points of the cavity

213

The heat transfer from the vertical wall can be obtained by the heat loss to the outside of the

214

cavity according to the energy balance of the vertical wall [33][34][35]. The calculation of heat

215

flux is based on stable conditions and the detailed calculation is as follows:

$$q_h = \frac{T_h - T_a}{A_w R_o} \quad (1)$$

216

Where  $T_h$  (K) is the temperature of inner vertical wall surface;  $T_a$  (K) is the temperature of the air

217

adjacent to the outer wall surface; and  $A_w$  ( $m^2$ ),  $R_o$  ( $K \cdot W^{-1}$ ) are the area of the vertical surface and

218

the thermal resistance between the inner surface and the outer ambient air respectively.  $R_o$  can be

219

measured as:

$$R_o = \frac{1}{(h_{o,rad} + h_{o,con})A_o} + \frac{L_w}{\lambda_w A_w} \quad (2)$$

220

where  $A_o$  ( $m^2$ ) is the area of the vertical wall that involves in heat exchange outside the cavity;  $\lambda_w$

221

( $W \cdot m^{-1} \cdot K^{-1}$ ) and  $L_w$  (m) are the thermal conductivity and the height of the vertical surface; and

222

$h_{o,rad}$  ( $W \cdot m^{-2} \cdot K^{-1}$ ),  $h_{o,con}$  ( $W \cdot m^{-2} \cdot K^{-1}$ ) are the radiant and convection heat transfer coefficient of the

223 outer vertical wall surface. Since the area of the outer wall surface is small enough with respect to  
224 the inner surface area of the entire air-conditioned room, the radiation heat transfer coefficient of  
225 the outer vertical wall surface can be calculated as:

$$h_{o,rad} = \frac{\sigma \varepsilon_o (T_{ow}^4 - T_c^4)}{(T_{ow} - T_c)} \quad (3)$$

226 where  $\sigma$  ( $\text{W}\cdot\text{m}^{-2}\cdot\text{K}^{-4}$ ) is the Stefan-Boltzmann constant,  $\varepsilon_o$  is the surface emissivity of the outer  
227 wall surface and is equal to 0.90,  $T_{ow}$  (K) and  $T_c$  (K) are the temperature of outer vertical wall  
228 surface and the inner surface of the air-conditioned room.

229 The convection heat transfer coefficient can be manipulated as:

$$h_{o,con} = \frac{Nu \lambda_w}{H_w} \quad (4)$$

230 where  $H_w$  (m) is the height of the vertical wall that involves in heat exchange outside the cavity,  
231  $Nu$  is the dimensionless parameter which is calculated using the correlation summarized in [36]  
232 based on the Rayleigh number.

233 The distributions of temperature points are shown as Fig. 3. Twenty-one RTDs were inserted  
234 at the following locations:  $x = 250$  mm as shown as Fig.1, at heights of 200 mm, 400 mm, 600  
235 mm away from the upper level of the heating unit along  $y$  axis, and at 0 mm, 8 mm, 15 mm, 30  
236 mm, 50 mm, 100 mm, 150 mm away from the vertical wall along the  $z$  axis. Six RTDs were  
237 mounted at  $x = 250$  mm and at heights ( $y$  direction) of 200 mm, 400 mm, and 600 mm; they were  
238 used to capture the outer vertical wall surface temperature and the temperature of the air adjacent  
239 to the outer surface. Four RTDs were used to measure the inner wall surface temperature of the  
240 air-conditioned room, and the temperature of the air-conditioned room was obtained by one RTD

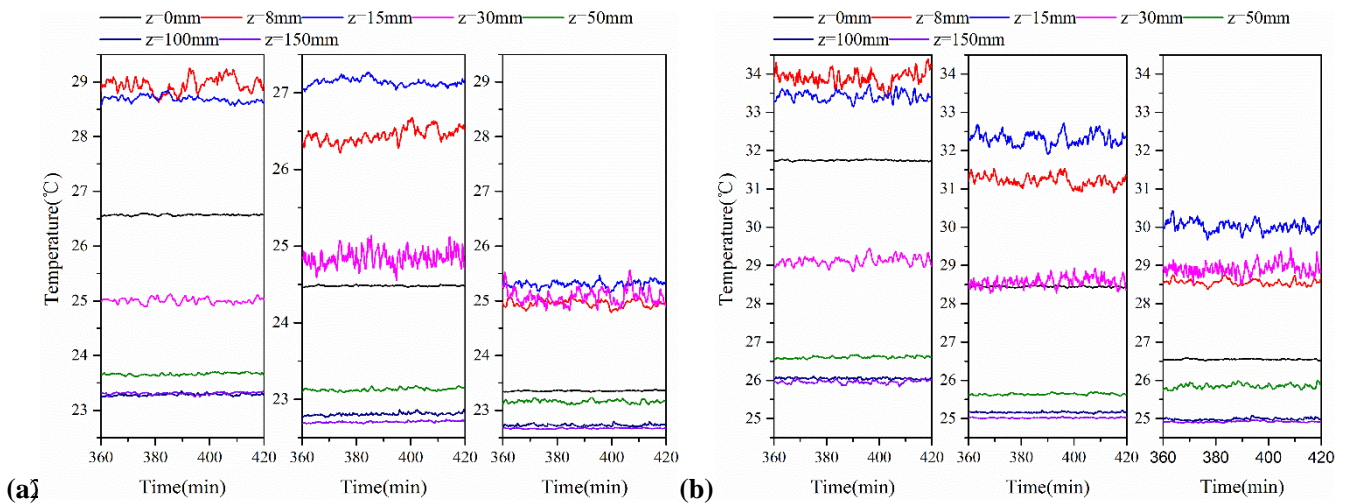
241 placed in the center of the room. The temperature series were stored in a computer by using the  
 242 Agilent 34980A multiplexer module. The sampling frequency was 0.1 Hz, with a response time of  
 243 10 s.

### 244 3. Results

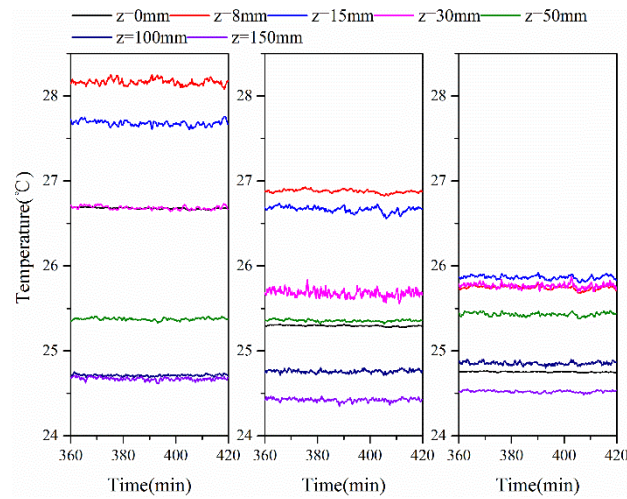
#### 245 3.1. Evaluation of temperature measurements

##### 246 3.1.1 Temperature distributions

247 Fig. 4 (a) - (c) show the temperature series of three cases in continuous time during the last  
 248 one hour. In Fig. 4 (a), (b), (c), the three subfigures from left to right represent the temperature  
 249 series at the heights of 200 mm, 400 mm and 600 mm. The temperature fluctuation near the  
 250 vertical surface is large at positions within 30 mm from the vertical wall. When the distance is  
 251 larger than 30 mm, the temperature tends to be relatively stable. As the heating temperature  
 252 increases, the air temperature near the vertical surface and the temperature fluctuation increases.  
 253 When mechanical ventilation is imposed, the near-wall air temperature and the temperature  
 254 fluctuation decreases.







(c)

**Fig. 4.** Temperature series of the  $y$ - positions, 200mm, 400mm, 600mm of the three cases (a) *Case 1* (b) *Case 2* (c) *Case 3*

256

257

258

259

The distributions of the time-averaged air temperature in the three cases presented in Fig. 5

260

show a similar distribution trend, which is an increase to a peak value and a gradual decrease. The

261

temperature of the heating unit and the ventilation condition not only affect the temperature

262

magnitude in the cavity, but also the temperature distribution characteristics. With an increase in

263

heating temperature, the air temperature in space increases significantly. The gradient of the air

264

temperature near the vertical surface increases dramatically as well. Under the effect of

265

mechanical ventilation, the air temperature of Case 3 tends to be more uniform, and the

266

temperature gradient near the vertical wall is largely weakened and approaches that of condition of

267

Case 1, even though the heating temperature is higher.

268

As stated in [3][15], due to the insufficiency of the thermal buoyancy driving force, part of

269

the hot airflow above the heat source deviates from the original flow direction. Due to flow

270

detachment, the air temperature distribution also exhibits the same detachment [3]. The hot air

271

flows upward for some distance until the buoyancy force is smaller than gravity and rushes to the

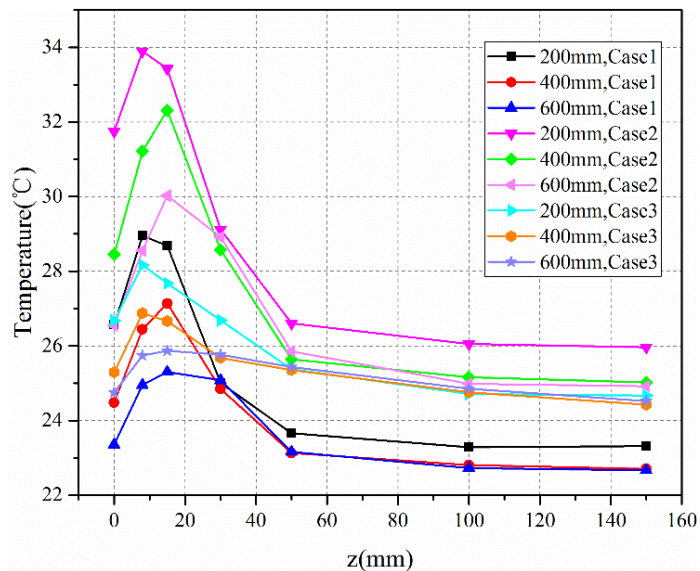
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inner regions of the cavity. At the height of 200 mm for all the cases, the air temperature increases

273

to a maximum value in the plane that is 8mm away from the vertical surface and then keeps

274 decreasing. If no temperature detachment exists, the trends at the other heights should also follow  
 275 the same law as that at the height of 200 mm. However, in Case 1 and Case 2, the maximum air  
 276 temperature occurs 15 mm away from the vertical surface at heights of 400 mm and 600 mm. In  
 277 consequence, there exists a temperature detachment. In Case 3, there also exists a temperature  
 278 detachment according to the changing tendency. Unlike the other natural cases, the maximum  
 279 temperature occurs in 8mm away from the vertical surface at the heights of 200 mm and 400 mm.  
 280 For the height of 600 mm, the maximum air temperature occurs 15mm away from the vertical wall.  
 281 The temperature detachment height rises to 600 mm under the effect of mechanical ventilation.

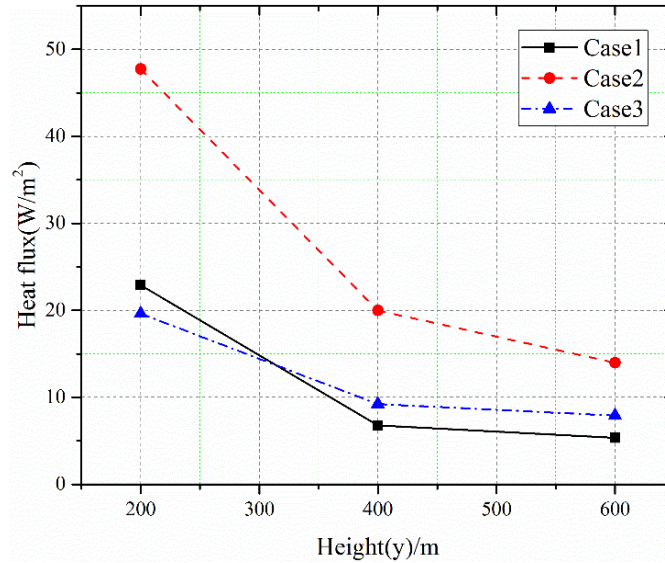


282  
 283 **Fig. 5.** The distribution of time-averaged temperature of three cases.

284 3.1.2 Heat transfer of the inner vertical wall

285 Based on the data processing method mentioned earlier, the inner heat transfer of the vertical  
 286 wall was calculated and is shown in Fig. 6. Similar to the trend of the inner surface temperature,  
 287 the heat flux of all the cases decreases with an increase in height. As the heating temperature  
 288 increases, the inner surface temperature of the vertical wall increases, and the potential difference  
 289 between the inner wall and the air outside the cavity increases; therefore, the heat loss to the outer

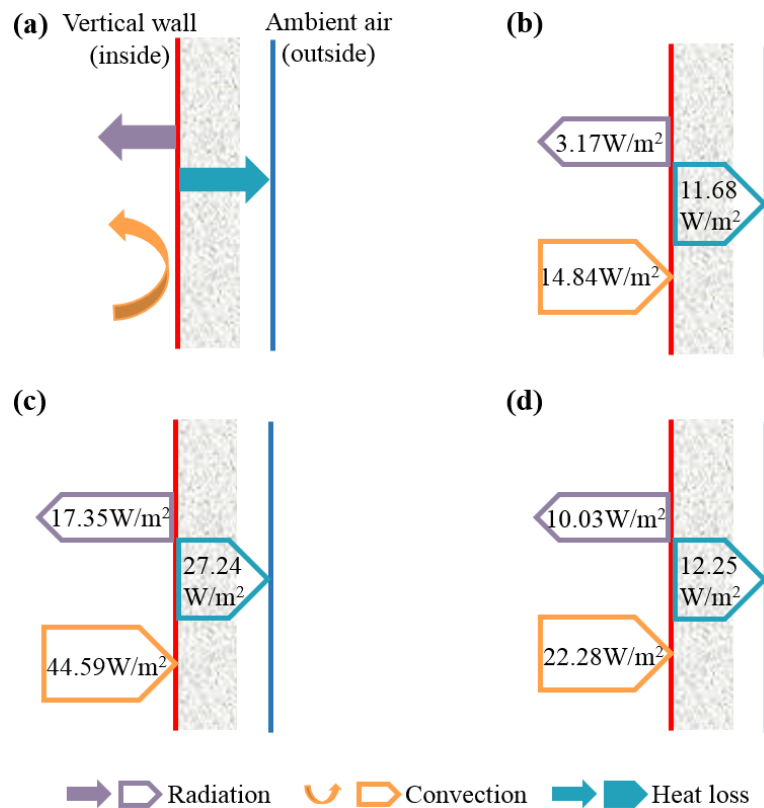
290 wall increases. Although the surface temperature of the heating unit in Case 3 is identical to that in  
 291 Case 2, the heat flux of the inner wall is much smaller than that of Case 2.



292  
 293 **Fig. 6.** The heat transfer of inner vertical wall

294 The energy balance of the vertical wall influences the heat transfer rate significantly. The  
 295 schematic heat balance of the vertical wall is shown as Fig. 7(a). In steady state, the heat storage  
 296 of the vertical wall is not taken into account. For the inner surface of the vertical wall, the heat  
 297 transfer process includes the radiative and convective heat transfer in the cavity and the heat loss  
 298 to the air adjacent to the outer wall. The radiative heat transfer rate of the inner wall is calculated  
 299 by establishing the radiative heat transfer network [37]. The other five inner surfaces are viewed as  
 300 a single surface since the temperature difference between each wall surface is not significant. The  
 301 single surface, the heating unit and the vertical wall constitute a closed radiation heat transfer  
 302 system. The surface thermal resistance is calculated by  $R_i = (1 - \varepsilon_i) / (A_i \varepsilon_i)$ , where  $\varepsilon$  is the surface  
 303 emissivity,  $A_i$  ( $m^2$ ) is the surface area. The space thermal resistance is calculated by  $R_{ij} = 1 / (A_i X_{ij})$ ,  
 304 where  $X_{ij}$  is the view factor from surface  $i$  to surface  $j$ . Since the vertical wall and the heating unit  
 305 are arranged in parallel, the view factor between the two surfaces is 0. According to heat balance,

306 the convective heat transfer rate is the sum of the heat loss and the radiative heat transfer rate. The  
 307 heat transfer rates of the three cases are shown in Fig. 7 (b)-(d). As shown in Fig.7, compared  
 308 with the radiative heat transfer, the convective heat transfer occupies a large proportion and has a  
 309 decisive role in the net heat flux of the vertical wall.



310

311 **Fig. 7.** (a) Schematic heat balance of the vertical wall, (b)-(d) detailed heat transfer rates of *Case 1*  
 312 to *Case 3*.

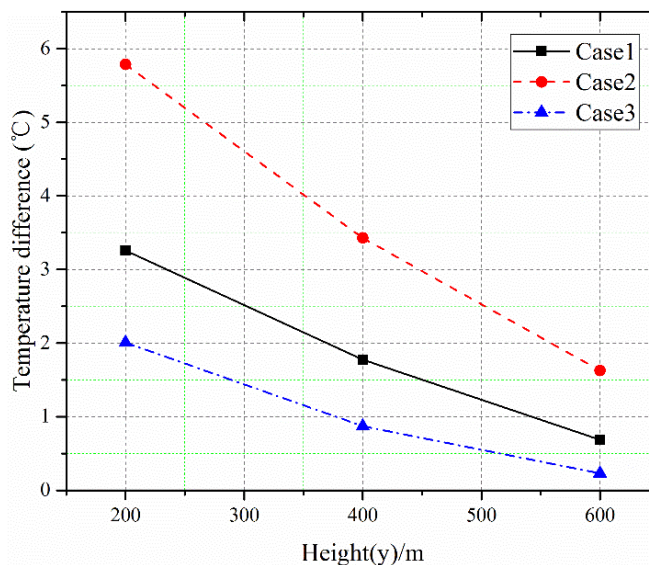
313 In the convective heat transfer process, the air flow rate affects the heat flux. Here, an  
 314 evaluation of the air flow rate based on the PIV measurement results was made. In the thin layer  
 315 near the wall, heat transfer occurs between the air and the wall and thus the air flow rate used to  
 316 calculate the heat flux should be close to the wall. We used the velocities in the plane that is 8mm  
 317 away from the vertical wall to calculate the air flow rate. The velocities in the 8mm plane are  
 318 averaged and the spatial-averaged velocities of Case 1 to Case 3 are 0.22m/s, 0.27m/s, and  
 319 0.96m/s. The air flow rates of Case 1 and Case 2 are shown to be smaller than that of Case 3.

320 Therefore, in this study, the impact of the air flow rate is not significant compared with the other  
321 influencing factors.

322 The thickness of the thermal boundary layer is an essential factor to evaluate heat transfer.  
323 With the increase of heating temperature, the heat transfer efficiency is enhanced, and the thermal  
324 boundary layer of Case 2 should be thinner than that of Case 1. As a result, through comparing the  
325 thickness of the thermal boundary layer of Case 2 and Case 3, the relative heat transfer efficiency  
326 of natural convection and the mixed convection can be determined. In this study, the thickness of  
327 the thermal boundary layer of Case 2 and Case 3 in the height of 200mm was compared. The  
328 Rayleigh number  $Ra (= g\beta(T_s - T_\infty)L^3 Pr/\nu^2)$  and the Reynolds number  $Re (= UL\nu^{-1})$  of Case  
329 2 and Case 3 are calculated, where  $\beta$  ( $K^{-1}$ ) is the thermal expansion coefficient,  $T_s$  (K) is the  
330 surface temperature,  $T_\infty$  (K) is the temperature of the mainstream and is regarded as the air  
331 temperature in the center of the cavity, and  $L$ (m) is the reference height of the vertical wall and is  
332 determined to be 0.2. Through calculation, it is found that the boundaries of the two cases are  
333 laminar. For the case of isothermal vertical surface under natural convection, the thickness of the  
334 thermal boundary layer can be calculated as  $\delta = 5.3\left(\frac{g\beta(T_s - T_\infty)}{\nu^2}\right)^{-0.25} L^{0.25}$  [36]. In this study,  
335 the vertical wall is non-isothermal. If  $T_s$  is taken as the temperature of the surface in the height of  
336 200mm, the calculated thickness of the boundary later would be larger than the real value. In  
337 consequence,  $T_s$  is taken as the average temperature of the surface in the region of 0-200mm. For  
338 the case of isothermal vertical surface under mixed convection, the thickness of the thermal  
339 boundary layer of the mixed convection can be calculated as  $\delta = \frac{5L}{Re^{0.5}}$  [36]. Then the  
340 thicknesses of the thermal boundary layers of Case 2 and Case 3 can be estimated to be 0.0125m

341 and 0.0066m. Therefore, the thickness of the thermal boundary layer of the natural convection  
342 cases are thicker than that of the mixed convection case. The change of heat flux is not consistent  
343 with the change of boundary layer, but presents an opposite trend.

344 However, as far as the experimental results in this study is concerned, the heat transfer cannot  
345 be only evaluated by the thickness of boundary layer. The thickness of the boundary layer is the  
346 combined effect of the flow and heat transfer. The basic factor that influencing the heat flux is the  
347 thermal and flow state of the air. In order to evaluate the factors contributing to the heat flux, the  
348 temperature and the velocity of the air should be inspected. In the previous response, we have  
349 analyzed the influence of the velocity and air flow rate. Next, the temperature difference between  
350 the near-wall air and the vertical wall is taken into accounted. Fig. 8 shows the temperature  
351 difference between the vertical wall and the air 100mm away from the vertical wall. As shown in  
352 Fig. 8, for Case 3, the temperature difference between the air and the vertical wall decreases  
353 significantly and even lower than that of Case 1. As a result, the temperature difference between  
354 the air and the vertical wall has a significant impact on the magnitude of the heat flux compared to  
355 the other factors.



356

357 **Fig. 8.** The temperature difference between the vertical wall and air 100mm away from the  
358 vertical wall

359 Furtherly, since the heat is transferred from the near-wall air to the vertical wall, the  
360 temperature of the near-wall air is critical to the magnitude of the heat flux of the vertical wall. In  
361 natural convection case, the air is heated and circulated in the cavity and no air enters in the air  
362 duct. Under the effect of the fan, the near-wall air of Case 3 flows out of the cavity along the outlet  
363 positioned in the top of the cavity. In this study, the air duct is not insulated and its heat  
364 conductivity coefficient is evidently larger than that of the cavity. A huge amount of the heat  
365 obtained in the cavity is released from the air duct to the air-conditioned room. With the  
366 continuous heat release process, the inlet air temperature of Case 3 is lowered significantly. As a  
367 result, the heat flux of Case 3 is similar to that of Case 1 even though the heating temperature is  
368 20°C higher.

### 369 **3.2. Evaluation of flow structures**

#### 370 3.2.1 Time averaged velocity distributions

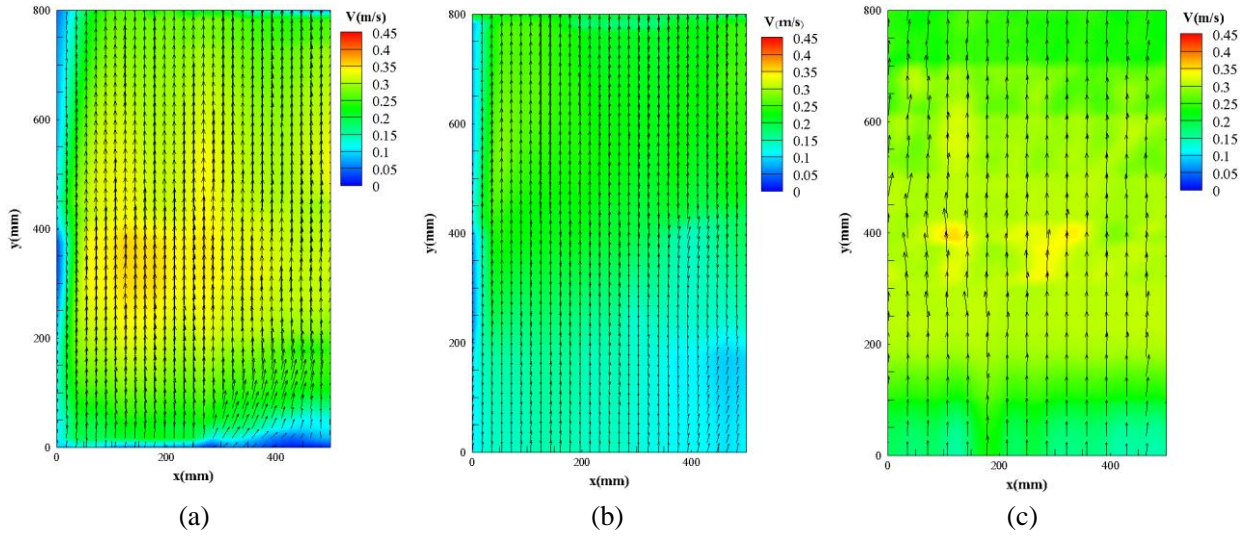
371 As mentioned previously, the two-dimensional velocity of the two planes that are 8 mm and  
372 15 mm away from the vertical wall were measured by PIV. For the convenience of illustration, the  
373 velocities of the plane that were 8 mm and 15 mm away from the vertical wall that were measured  
374 by PIV are denoted as PIV8 and PIV15, respectively. The measured results of thermal image  
375 velocimetry are denoted as TIV. By averaging the instantaneous velocities, the time-averaged  
376 velocity fields measured by PIV and TIV under three cases can be obtained (Fig. 9 to Fig. 11).

377 Fig. 9 shows the time-averaged velocity field of Case 1. From the perspective of the spatial  
378 distribution, the measured velocity shows vertical variation to some extent. As the height increases,

379 the velocity gradually increases and reaches at maximum value at a height of 200 mm. The  
380 velocity is maintained at a high level until the height increases to 400 mm. As the height continues  
381 to increase, the velocity decreases gradually due to an insufficient buoyancy force. The  
382 distribution trend of Case 2 is similar to that of Case 1. However, due to an increase in heating  
383 temperature, the velocity distribution exhibits a more pronounced spatial difference in the vertical  
384 direction. In Case 2, the peak velocity appears at a height of 400-600 mm, which is higher than  
385 that of Case 1. Correspondingly, the maximum velocity value increases as well. Different from the  
386 distributions presented in Case 1 and Case 2, the velocity distributions of Case 3 show a  
387 decreasing trend with an increase in height. In Case 3, natural convection and mechanical  
388 ventilation have a synergistic effect on velocity distribution. As can be seen from Fig. 10, the  
389 influence of mechanical ventilation is stronger than that of natural convection and the velocity  
390 distribution characteristics are mainly determined by mechanical ventilation.

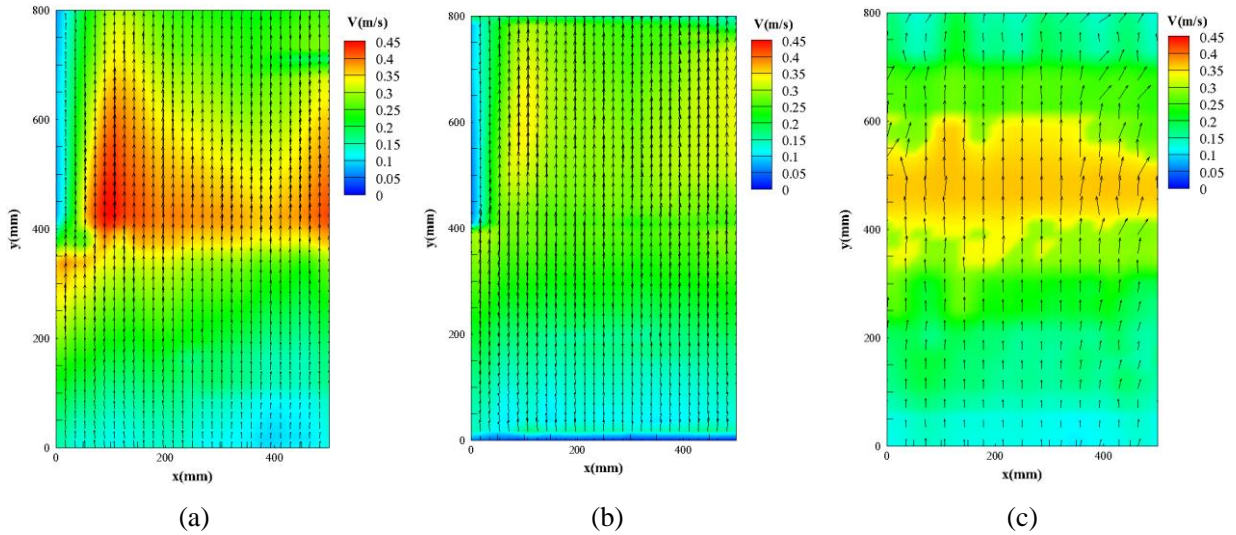
391 Further inspection of the relative magnitude of PIV15 and PIV8 shows that PIV15 is smaller  
392 than PIV8. The plume originates from the heating unit and rises to the upper quiescent region. As  
393 the perpendicular distance from the heating unit increases, the scale of the plume increases,  
394 eventually dissipating in the quiescent region owing to a decrease in buoyancy force caused by the  
395 cooling of cold air and viscous effects [36]. Due to an insufficient thermal buoyancy, the velocity  
396 decreases with an increase in perpendicular distance. As a result, the velocity in the plane that is  
397 15 mm away from the vertical wall is lower than the velocity in the plane that is 8mm away from  
398 the vertical wall.





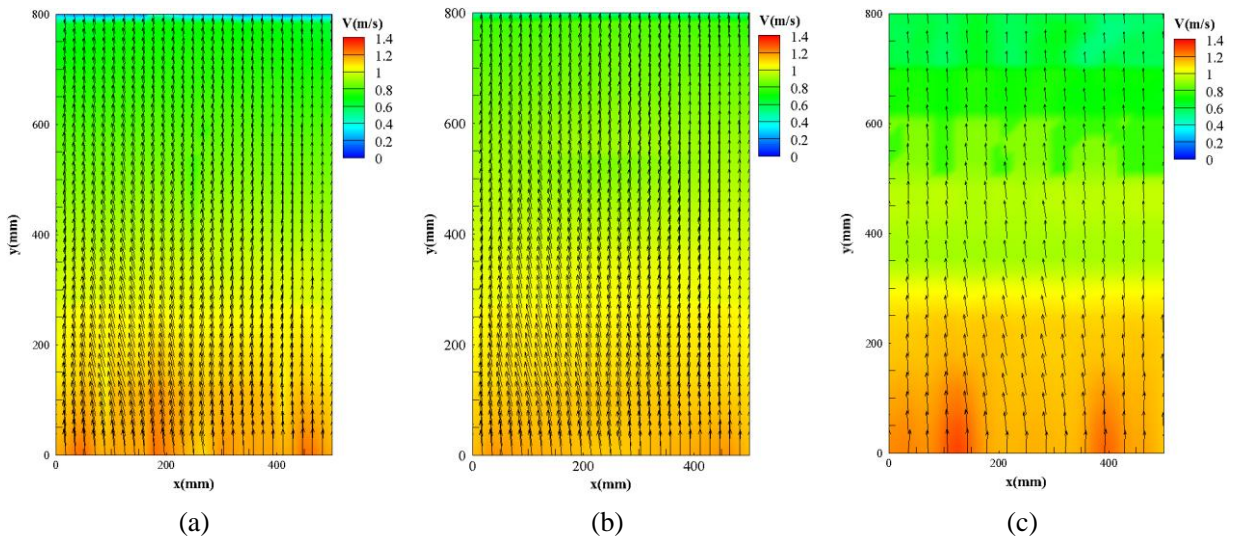
399 **Fig. 9.** Time-averaged velocity of *Case 1*, (a) PIV8, (b) PIV15, (c) TIV

400



401 **Fig. 10.** Time-averaged velocity of *Case 2*, (a) PIV8, (b) PIV15, (c) TIV

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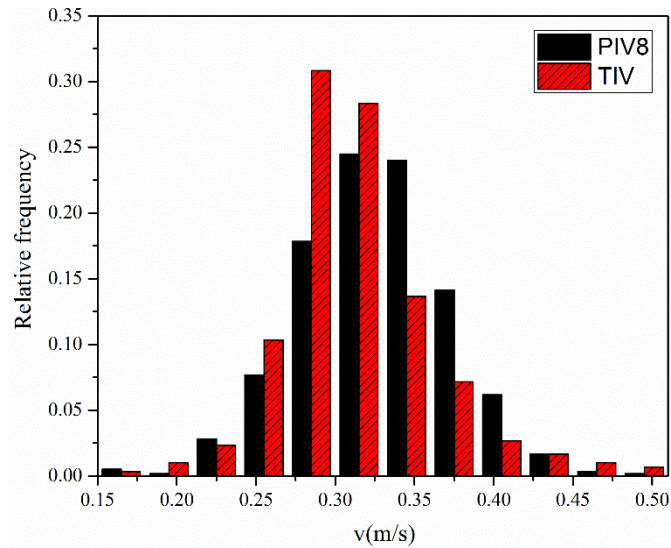
402 **Fig. 11.** Time-averaged velocity of *Case 3*, (a) PIV8, (b) PIV15, (c) TIV

403 On the whole, the velocity fields measured by TIV are close to those measured by PIV,  
404 especially the measurement results of PIV8. However, in the three Cases, differences between TIV  
405 and PIV8 can be observed. The differences are directly related to the difference in measurement  
406 method principles. In the interaction between the wall surface and flow structures, heat is  
407 exchanged with the ejections and sweep of intermittent eddies. An essential factor that used to  
408 evaluate the heat transfer is the thickness of the boundary layer. As the heating temperature  
409 increases or mechanical ventilation is imposed, the thickness of boundary layer becomes thinner  
410 and the heat transfer is more intense. In TIV, the velocity is obtained by detecting the fluctuations  
411 of surface thermal spots. The temperature pattern captured by the infrared camera is closely  
412 related to the boundary layer [38]. As the thickness of the boundary layer changes, the “position”  
413 of the flow field measured by TIV changes accordingly. However, in the PIV experiment, the flow  
414 field of the planes that are 8 mm and 15 mm away from the vertical surface were measured, and  
415 thus, the measurement position is always fixed. With the change in the thickness of the boundary  
416 layer, the relative positions of PIV8 and TIV changes, and thus, the measurement results of TIV  
417 and PIV8 exhibit a certain difference in the three cases.

418 Overall, in natural convection and mixed convection, TIV and PIV8 are close, and both are  
419 more consistent in magnitude and trend. The position where the maximum velocity occurs and the  
420 magnitude of velocity can be accurately interpreted by TIV. Therefore, TIV has a certain  
421 feasibility in reflecting the spatial distribution characteristics of the near-wall airflow.

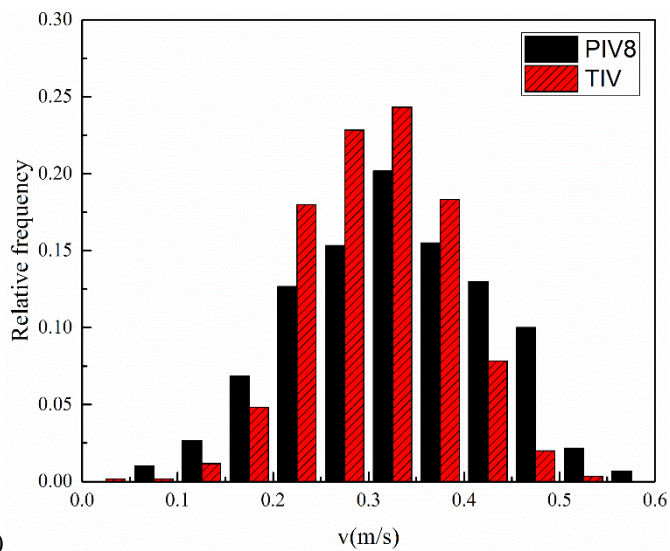
422 The histograms of velocities at  $x = 250$  mm and  $y = 200$  mm for all the cases are shown in  
423 Fig. 12. As a whole, the distributions of PIV8 and TIV are similar, which shows that the  
424 statistical results are reliable. It is can be observed that, in PIV8, the large velocity values account

425 for a larger proportion than that in the TIV for the three cases. This is because that there exists a  
 426 lag between the surface temperature fluctuation and the near-wall velocity fluctuation. When the  
 427 near-wall airflow changes, the surface shows the trace of the airflow in the form of surface  
 428 temperature fluctuations. However, due to the heat storage of the vertical wall, the surface  
 429 temperature cannot change immediately when the near-wall flow changes. As a result, the surface  
 430 temperature fluctuations are weaker than those in the ideal situation, and the TIV measurement  
 431 results tend to be distributed in the low velocity range.



432

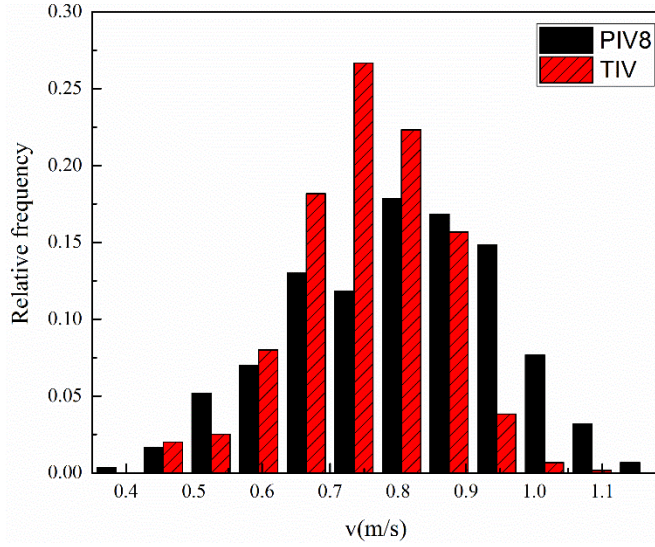
(a)



433

(b)





(c)

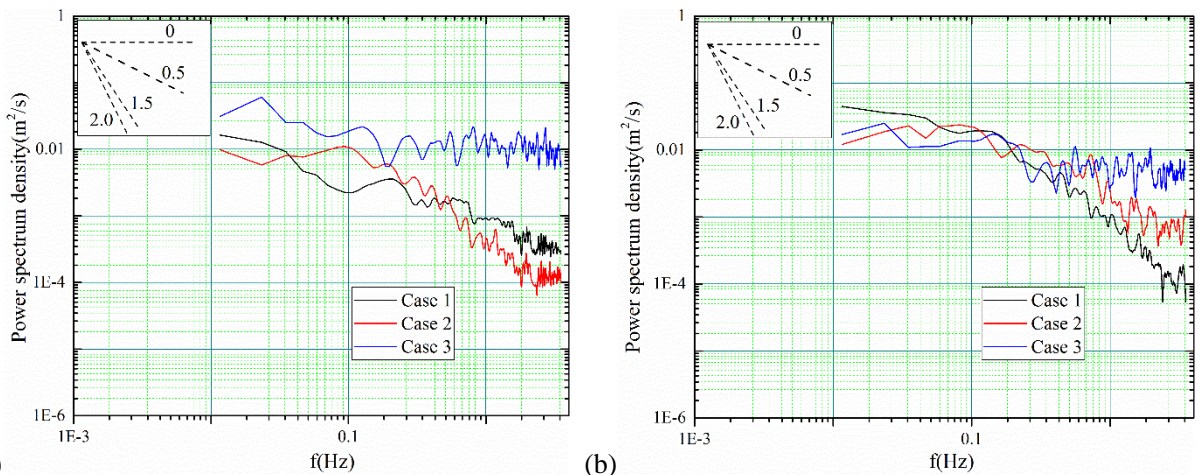
**Fig. 12.** The histograms of velocity at  $x=250\text{mm}$ ,  $y=200\text{mm}$ , (a) *Case 1*, (b) *Case 2*, (c) *Case 3*

### 3.2.2 Power spectrum analysis

The power spectrum analysis of the airflow fluctuation can be used to reveal the energy distributions in a frequency range and characterize the turbulent airflow [39]. The fluctuating velocity ( $v'$ ) is obtained by subtracting the instantaneous velocity from the time-averaged velocity. Based on the velocity fluctuation, the power spectrum density function can be defined as follows:

$$\int_0^{\infty} E(f)df = \overline{v'^2} \quad (5)$$

where  $E(f)$  ( $\text{m}^2/\text{s}$ ) is the power spectrum density function,  $f$  is the frequency (Hz).

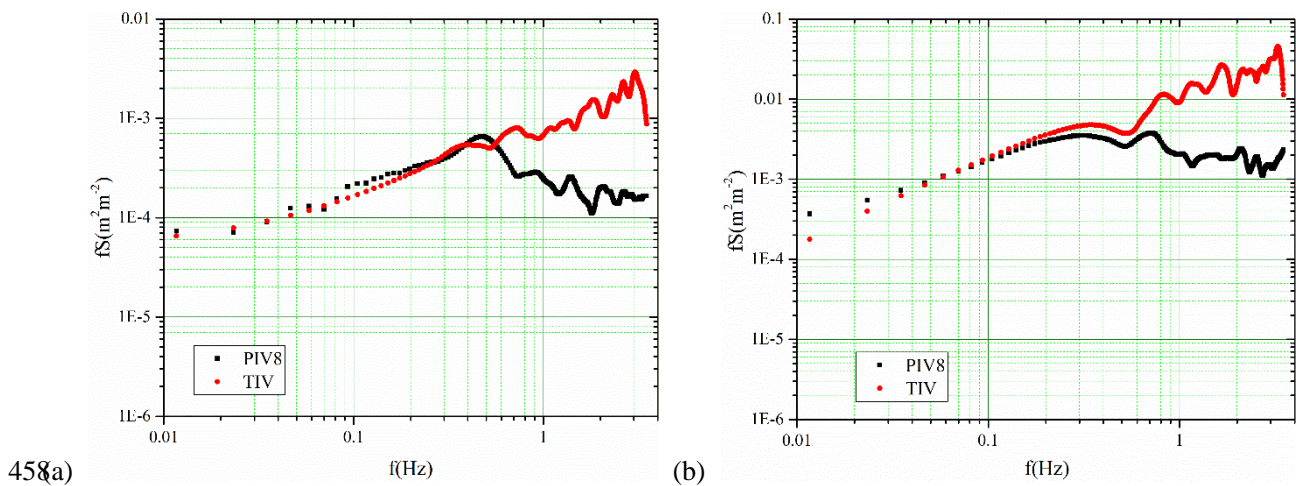


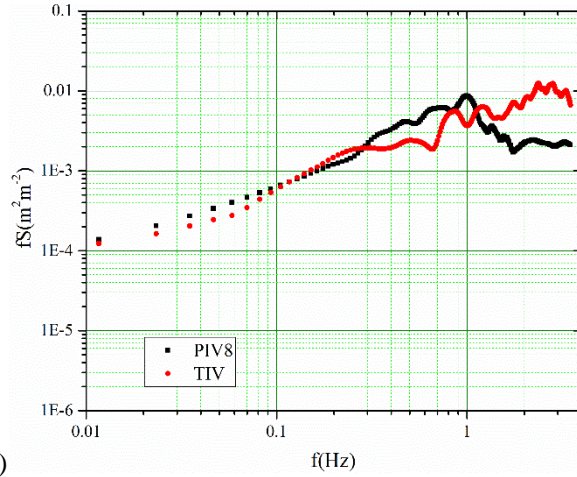
**Fig. 13.** Power spectra (a) at the height of 200mm; (b) at the height of 600mm

Fig. 13. (a) and (b) show the power spectrum density measured by PIV at the heights of 200

445 mm and 600 mm, respectively. The spectrum distributions of the velocity fluctuation of the three  
 446 cases show a similar tendency, but there are still some differences in the inertial subrange. The  
 447 negative slope of the logarithmic power spectrum curve of natural convection are between 1.5-2.0,  
 448 while the negative slope of the mixed convection case is within 0-0.5. This phenomenon was also  
 449 observed by [40]. The negative slope of the inertial subrange in natural convection are close to 5/3,  
 450 showing that the flows are fully developed. The turbulence of Case 3 differs from that in the other  
 451 cases, as it is influenced by the jet at the bottom of the cavity. With the diffusion of the jet airflow  
 452 in the cavity, the average velocity decreased slightly, and the turbulence decreased slightly  
 453 accordingly.

454 The pre-multiplied power spectrum is obtained by multiplying the power spectrum of  
 455 velocity and the wave number, which can be used to show the contribution of different wave  
 456 number on kinetic energy. The pre-multiplied velocity spectra at  $x = 250$  mm,  $y = 200$  mm of TIV  
 457 and PIV8 of the three cases are shown as Fig. 14.





(c)

**Fig. 14.** Pre-multiplied power spectra of (a) *Case 1*, (b) *Case 2*, (c) *Case 3*

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As can be seen from Fig. 14, the spectral shapes of the three cases all fit well in the low-frequency regions. As stated in [41], the turbulent flow consists of large scale and small scale structures. The large scale structures fluctuate with low frequency and represent the main motion of turbulent flow. The small scale structures are distributed in the high-frequency region and are advected by large scale structures. The great fitness in the low-frequency region means that the measurement results of TIV can reflect the main flow characteristics of turbulent flow motion. In the higher frequency regions, the spectral curves separate and the pre-multiplied power spectrum of TIV is more energetic than PIV, which is also observed in [25]. Due to the wall shear stresses, the closer to the wall, the more energetic the flow fluctuation is. The peak of the pre-multiplied spectra appears in the higher frequency region, which means that the airflow field evaluated by TIV may be closer to the vertical wall than the airflow field measured by PIV.

#### 4. Discussion

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In the foregoing section, we mentioned that the measurement results of TIV are influenced by the boundary layer. By analyzing the pre-multiplied spectrum of PIV and TIV, it is deduced that the position of the airflow field evaluated by TIV may be closer to the wall than the airflow field

476 of the fixed plane measured by PIV. Although the measurement position of TIV cannot be directly  
477 determined, the dynamic characteristics of the PIV and TIV measurements provide a new  
478 perspective for estimating the information of the near-wall airflow field. We attempted to  
479 indirectly judge whether the result of the TIV measurement is affected by the boundary layer by  
480 judging the relative position of the two airflow fields measured by PIV and TIV and the boundary  
481 layer.

482 From the perspective of physical structure, turbulent flow can be regarded as a superposition  
483 of vortices with various sizes. Large scale vortices fluctuate with low frequency and extract energy  
484 from the mainstream. Through the interaction of vortices, energy is gradually transferred to the  
485 small-scale vortices which fluctuate with high frequency. Finally, due to fluid viscosity, small  
486 scale vortices disappear and the mechanical energy is converted into thermal energy of the fluid.  
487 The spectrum of temperature can be calculated as follows:

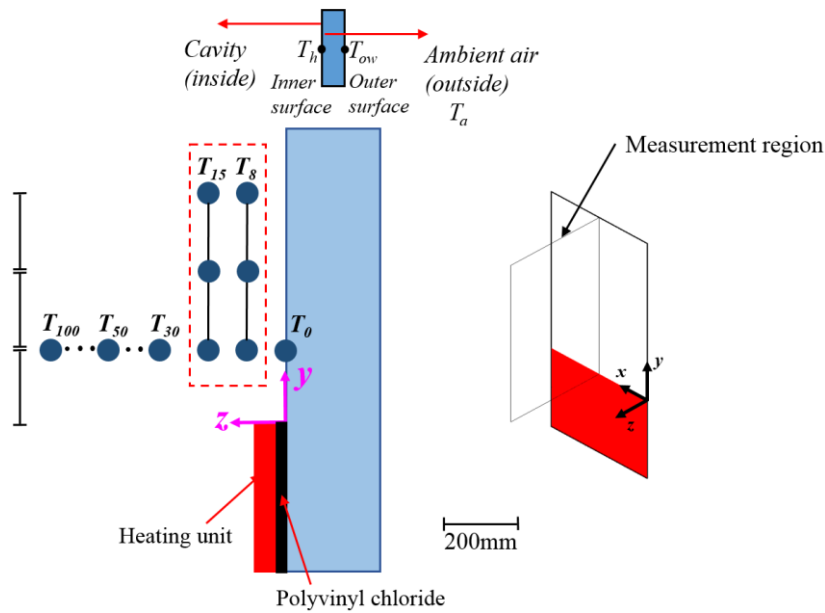
$$E_{\theta} = \int_0^{\infty} E_{\theta}(k) dk \quad (6)$$

488 where  $E_{\theta}$  is the power spectrum density function of temperature,  $k$  is the wave number.

489 The exponential slope of the region between the flat low-frequency and the sharp drop at the  
490 high frequency represents the efficiency of energy transformation from the large scale to the small  
491 scale structures. In Zhang' study [42], it was found that the efficiency of energy transformation  
492 near the vertical surface varies greatly and reaches a maximum value near the boundary layer.  
493 Outside the boundary layer, the energy transformation efficiency almost remains constant with a  
494 constant exponent. As a result, it appears that the efficiency of the transformation from the large  
495 scale structure to the small scale structures can be used to estimate the position of the boundary  
496 layer. Based on this characteristic, we measured the temperature fluctuation of the air near the

497 vertical surface and explored the slope of the power law distribution region.

498 An appropriate temperature sampling frequency is required to obtain the real-time  
 499 information about the temperature field. However, due to the limitations of the measurement  
 500 instrument, the sampling frequency is always restricted. The Kolmogorov time scale is used to  
 501 evaluate whether the measurement frequency is sufficient. The Kolmogorov time scale can be  
 502 calculated as:  $\tau=(\nu/\varepsilon)^{0.5}$ ,  $\varepsilon=U^3/H$ , where  $\nu$  is the turbulent viscosity,  $\varepsilon$  is the turbulent  
 503 dissipation rate,  $U$  is root mean square velocity,  $H$  is the turbulence length scale. In this study, the  
 504 Kolmogorov time scale is 0.245 s for Case 1. The temperature fluctuations were collected by  
 505 Angilent 34980A with a frequency of 7 Hz. The response time is 0.143 s, which is smaller than the  
 506 Kolmogorov time scale for Case 1. In consequence, the data collected at 7 Hz can be used to  
 507 evaluate the energy transformation efficiency.



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**Fig. 15.** Schematic distribution of temperature measurement points

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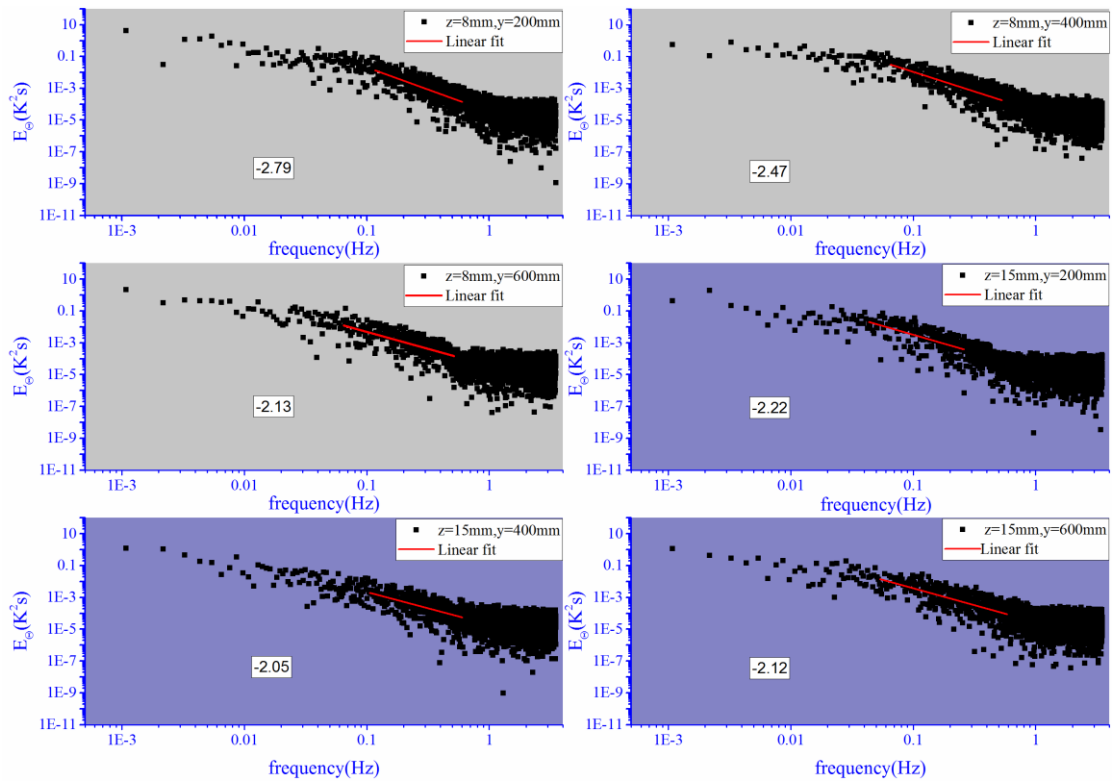
512

As shown in Fig. 15, the temperature fluctuations of the ten points were measured. The power spectrum distributions of the six points in the planes that are 8 mm and 15 mm away from the vertical surface are shown in Fig. 16. The six points have the same arrangement as the former



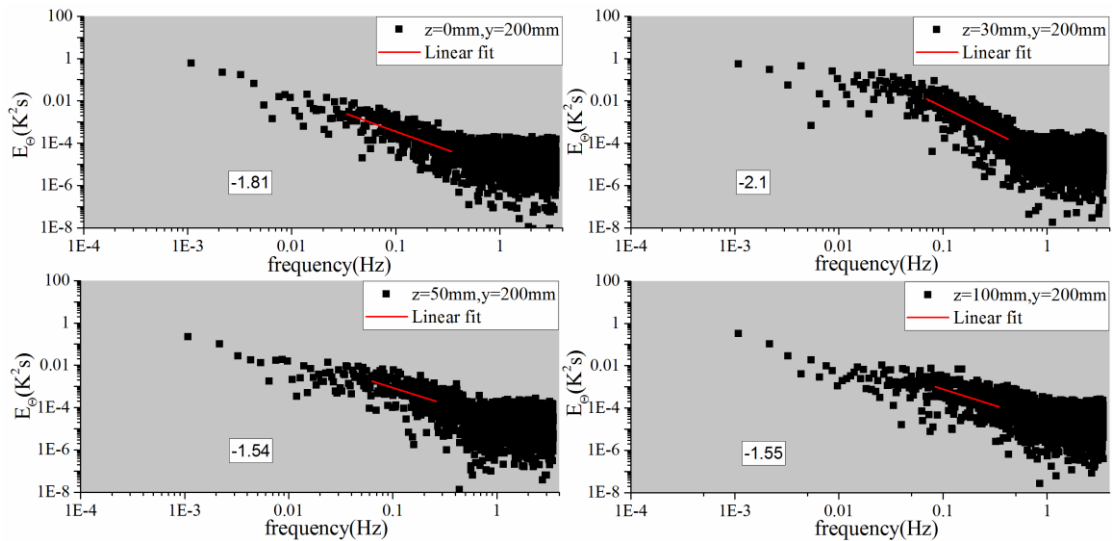
513 temperature measurement arrangement. The power spectrum distributions of the other four points  
514 located at the height of 200 mm, and are 0 mm, 30 mm, 50 mm, 100 mm away from the vertical  
515 surface are presented in Fig. 17. The line are fitted in the power law distribution region and the  
516 slopes are shown in figures.

517 For the plane that is 8 mm away from the vertical surface, energy transfer efficiency  
518 decreases with an increase in height. As for the plane that is 15 mm away from the vertical surface,  
519 heat transfer efficiency is almost constant and is smaller than that for the plane that is 8 mm away  
520 from the vertical surface. Combining Fig. 16 and Fig. 17, it can be found that the energy  
521 transformation efficiency increases from 0 mm to 8 mm. The efficiency decreases from 8 mm to  
522 15 mm and remains almost constant within 15 mm to 30 mm from the vertical surface. When the  
523 distance is increased to 50 mm, the efficiency decreases greatly. When distance increases and for  
524 to the regions far away from the vertical surface, the efficiency reduces due to a small temperature  
525 gradient and weak convection. According to the aforementioned analysis, the efficiency increased  
526 to its maximum value near the boundary layer and remains almost constant outside the boundary  
527 layer. It can, therefore, be deduced that the plane that is 8 mm away from the vertical surface is  
528 near the boundary layer, and the plane that is 15 mm away from the vertical surface is outside the  
529 boundary layer. In combination with the fact that the TIV velocity is close to that of PIV8, it can  
530 be identified that the velocity plane measured by TIV is also near the boundary layer. Although  
531 unable to establish a direct relation, it is ascertained that the measured results of TIV can reflect  
532 the dynamic characteristics of the near-wall region in either natural or mixed convection.



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**Fig. 16.** Power spectrum of temperature of the points in the plane of 8mm and 15 mm away from the vertical surface



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538

**Fig. 17.** Power spectrum of temperature of the points that are 0mm, 30mm, 50mm and 100mm away from the vertical surface

539 Although the dimension and shape of the heating unit in this study are somewhat different  
540 from those of actual radiator, the heat transfer process of the heating unit is similar to that of an  
541 actual radiator under full scale room conditions. The influence of the surface temperature of the  
542 heating unit and the mechanical ventilation on the thermodynamic state of the near-wall airflow

543 over the heating unit found in this study can be referred in the research of actual heating room. In  
544 addition, TIV is proved to be feasible in the measurement of the near-wall airflow over a heating  
545 unit and can be used as an alternative indoor velocity measurement method to make up for the fact  
546 that it is difficult to use PIV in full scale measurements. Based on the measurement and analyzing  
547 method, future work can be carried out to investigate the thermodynamic state of the airflow over  
548 a radiator in a full scale room environment.

## 549 **5. Conclusions**

550 This paper analyzed the thermodynamic state of the near-wall airflow over an idealized  
551 heating unit. The air temperature was measured by RTDs, and the near-wall velocity distributions  
552 were measured by the PIV method and TIV method. By analyzing the temperature and velocity  
553 measurement results, the following conclusions can be drawn:

554 (a) Overall, TIV and PIV8 are close and consistent in magnitude and trend. TIV has certain  
555 feasibility in reflecting the spatial distribution characteristics of the near-wall airflow in  
556 both natural convection and mixed convection. The spectral shapes of TIV and PIV fit  
557 well in low-frequency regions. The great fitness in the low-frequency region means that  
558 the measurement results of TIV can reflect the main flow characteristics of turbulent  
559 flow motion. TIV can be used to obtain the near-wall velocity according to the wall  
560 surface temperature fluctuations and the obtained velocity is closely associated with the  
561 boundary layer. Due to the heat storage of the vertical wall, lag exists between the  
562 surface temperature fluctuation and the near-wall velocity fluctuation.

563 (b) Influenced by thermal buoyancy, the near-wall velocity over the heating unit in natural

564 convection case shows vertical variation. With an increase in heating temperature, the  
565 velocity distribution exhibits a more pronounced spatial difference in the vertical  
566 direction, and the position at which the maximum velocity occurs moves upward.  
567 Correspondingly, the value of the maximum velocity increases with an increase in  
568 heating temperature as well. The flow under natural convection cases is in fully  
569 developed turbulent, and the flow state of mixed convection is influenced by the  
570 characteristics of supply air. The influence of mechanical ventilation is far stronger than  
571 that of natural convection in mixed convection case, and as the vertical distance from the  
572 air inlet increases, the air velocity gradually decreases.

573 (c) The temperature fluctuation near the vertical surface is large at locations within 30 mm  
574 away from the vertical surface. The temperature of the air near the vertical wall has a  
575 significant influence on the heat transfer rate of the vertical wall. With an increase in  
576 heating temperature, the air and surface temperature gradients increase, especially those  
577 near the vertical surface. Accordingly, the near-wall heat transfer rate of the vertical wall  
578 increases as well. Under the effect of mechanical ventilation, the air and surface  
579 temperatures decrease and tend to be more uniform, and the temperature gradient is  
580 reduced. The heat transfer rate of the vertical wall decreases. Due to an insufficient of  
581 driving force, there exists a temperature detachment over the heating unit. Influenced by  
582 mechanical ventilation, the detachment position of the mixed convection case is higher  
583 than those of the natural convection cases.

#### 584 **Acknowledgements**

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## 587 **References**

- 588 [1] J.A. Myhren, S. Holmberg, Performance evaluation of ventilation radiators, *Appl. Therm.*  
589 *Eng.* 51 (2013) 315–324. doi:10.1016/j.applthermaleng.2012.08.030.
- 590 [2] J.A. Myhren, S. Holmberg, Design considerations with ventilation-radiators: Comparisons to  
591 traditional two-panel radiators, 41 (2009) 92–100. doi:10.1016/j.enbuild.2008.07.014.
- 592 [3] G. Sevilgen, M. Kilic, Numerical analysis of air flow , heat transfer , moisture transport and  
593 thermal comfort in a room heated by two-panel radiators, 43 (2011) 137–146.  
594 doi:10.1016/j.enbuild.2010.08.034.
- 595 [4] A. Jahanbin, E. Zanchini, Effects of position and temperature-gradient direction on the  
596 performance of a thin plane radiator, *Appl. Therm. Eng.* 105 (2016) 467–473.  
597 doi:10.1016/j.applthermaleng.2016.03.018.
- 598 [5] W. Zhao, Y. Hu, Y. Wang, W. Qin, Thermal performance of a suspended ceiling fi n heat  
599 transfer panel with drain pan, *Build. Environ.* 144 (2018) 622–630.  
600 doi:10.1016/j.buildenv.2018.09.004.
- 601 [6] Y. Lyu, X. Wu, C. Li, H. Su, L. He, Numerical analysis on the effectiveness of warm water  
602 supply in water flow window for room heating, *Sol. Energy.* 177 (2019) 347–354.  
603 doi:10.1016/j.solener.2018.11.033.
- 604 [7] S. Xu, R. Ding, J. Niu, G. Ma, Investigation of air-source heat pump using heat pipes as heat  
605 radiator, *Int. J. Refrig.* 90 (2018) 91–98. doi:10.1016/j.ijrefrig.2018.03.025.
- 606 [8] H.B. Awbi, A. Hatton, Natural convection from heated room surfaces, *Energy Build.* 30  
607 (1999) 233–244.
- 608 [9] T. Cholewa, R. Anasiewicz, A. Siuta-Olcha, M.A. Skwarczynski, On the heat transfer  
609 coefficients between heated/cooled radiant ceiling and room, *Appl. Therm. Eng.* 117 (2017)  
610 76–84. doi:10.1016/j.applthermaleng.2017.02.019.
- 611 [10] Y. Sun, Y. Zhang, An Overview of Room Air Motion Measurement: Technology and  
612 Application, *HVAC&R Res.* 13 (2007) 929–950. doi:10.1080/10789669.2007.10391463.
- 613 [11] S. Fu, P.H. Biwale, C. Mathis, Particle tracking velocimetry for indoor airflow field: A review,  
614 *Build. Environ.* 87 (2015) 34–44. doi:10.1016/j.buildenv.2015.01.014.
- 615 [12] C. Cierpka, R. Hain, N.A. Buchmann, Flow visualization by mobile phone cameras, *Exp.*  
616 *Fluids.* 57 (2016). doi:10.1007/s00348-016-2192-y.
- 617 [13] A. Sattari, Particle image velocimetry visualization and measurement of airflow over a  
618 wall-mounted radiator, *Int. J. Vent.* 14 (2015) 289–302.  
619 doi:10.1080/14733315.2015.11684087.
- 620 [14] W. Du, Y. Liu, H. Yuan, S. Qiao, S. Tan, Experimental investigation on natural convection  
621 and thermal stratification of IRWST using PIV measurement, *Int. J. Heat Mass Transf.* 136  
622 (2019) 128–145. doi:10.1016/j.ijheatmasstransfer.2019.01.067.
- 623 [15] Q. Wu, C. an Zhu, L. Liu, J. Liu, Z. Luo, Two-dimensional flow visualization and velocity  
624 measurement in natural convection near indoor heated surfaces using a thermal image

- 625 velocimetry method, *Appl. Therm. Eng.* 146 (2019) 556–568.  
 626 doi:10.1016/j.applthermaleng.2018.10.023.
- 627 [16] W. Gao, R.H. Shaw, K.T. Paw U, Observation of organized structure in turbulent flow within  
 628 and above a forest canopy, *Boundary-Layer Meteorol.* 47 (1989) 349–377.  
 629 doi:10.1007/BF00122339.
- 630 [17] F. Castellví, C. Cammalleri, G. Ciraolo, F. Rossi, Daytime sensible heat flux estimation over  
 631 heterogeneous surfaces using multitemporal land-surface temperature observations. *Water*  
 632 *Resour. Res.* 52 (2016) 3457–3476.
- 633 [18] G.M. Carlomagno, G. Cardone, Infrared thermography for convective heat transfer  
 634 measurements, 2010. doi:10.1007/s00348-010-0912-2.
- 635 [19] M. Aminzadeh, D. Breitenstein, D. Or, Characteristics of Turbulent Airflow Deduced from  
 636 Rapid Surface Thermal Fluctuations: An Infrared Surface Anemometer, *Boundary-Layer*  
 637 *Meteorol.* (2017) 1–16. doi:10.1007/s10546-017-0279-5.
- 638 [20] E. Haghighi, D. Or, Thermal signatures of turbulent airflows interacting with evaporating thin  
 639 porous surfaces, *Int. J. Heat Mass Transf.* 87 (2015) 429–446.
- 640 [21] C.J. Legleiter, P.J. Kinzel, J.M. Nelson, Remote measurement of river discharge using  
 641 thermal particle image velocimetry (PIV) and various sources of bathymetric information, *J.*  
 642 *Hydrol.* 554 (2017) 490–506. doi:10.1016/j.jhydrol.2017.09.004.
- 643 [22] T. Lopez, H.E. Thomas, A.J. Prata, A. Amigo, D. Fee, D. Moriano, Volcanic plume  
 644 characteristics determined using an infrared imaging camera, *J. Volcanol. Geotherm. Res.* 300  
 645 (2014) 148–166. doi:10.1016/j.jvolgeores.2014.12.009.
- 646 [23] A. Tiddens, K. Risthaus, M. Röger, H. Stadler, B. Hoffschmidt, Induced Infrared  
 647 Thermography: Flow visualizations under the extreme conditions of an open volumetric  
 648 receiver of a solar tower, *Int. J. Heat Fluid Flow.* 65 (2017) 105–113.  
 649 doi:10.1016/j.ijheatfluidflow.2017.04.002.
- 650 [24] Z.C. Huang, K.S. Hwang, Measurements of surface thermal structure, kinematics, and  
 651 turbulence of a large-scale solitary breaking wave using infrared imaging techniques, *Coast.*  
 652 *Eng.* 96 (2015) 132–147. doi:10.1016/j.coastaleng.2014.12.005.
- 653 [25] A. Inagaki, M. Kanda, S. Onomura, H. Kumemura, Thermal Image Velocimetry,  
 654 *Boundary-Layer Meteorol.* 149 (2013) 1–18. doi:10.1007/s10546-013-9832-z.
- 655 [26] Y. Fan, Y. Li, S. Yin, Non-uniform ground-level wind patterns in a heat dome over a  
 656 uniformly heated non-circular city, *Int. J. Heat Mass Transf.* 124 (2018) 233–246.  
 657 doi:10.1016/j.ijheatmasstransfer.2018.03.069.
- 658 [27] E. Lev, M. Spiegelman, R.J. Wysocki, J.A. Karson, Investigating lava flow rheology using  
 659 video analysis and numerical flow models, *J. Volcanol. Geotherm. Res.* 247–248 (2012)  
 660 62–73. doi:10.1016/j.jvolgeores.2012.08.002.
- 661 [28] A. Garai, J. Kleissl, Air and Surface Temperature Coupling in the Convective Atmospheric  
 662 Boundary Layer, *J. Atmos. Sci.* 68 (2011) 2945–2954. doi:10.1175/JAS-D-11-057.1.
- 663 [29] M. Katurji, P. Zawar-reza, Forward-Looking Infrared Cameras for Micrometeorological  
 664 Applications within Vineyards, *Sensors.* 16 (2016) 2–11. doi:10.3390/s16091518.
- 665 [30] A.M. Grudzielanek, J. Cermak, Capturing Cold-Air Flow Using Thermal Imaging,  
 666 *Boundary-Layer Meteorol.* 157 (2015) 321–332. doi:10.1007/s10546-015-0042-8.
- 667 [31] A. Koca, G. Cetin, Experimental investigation on the heat transfer coefficients of radiant  
 668 heating systems: Wall, ceiling and wall-ceiling integration, *Energy Build.* 148 (2017)

- 669 311–326. doi:10.1016/j.enbuild.2017.05.027.
- 670 [32] X. Zhang, G. Su, J. Yu, Z. Yao, F. He, PIV measurement and simulation of turbulent thermal  
671 free convection over a small heat source in a large enclosed cavity, *Build. Environ.* 90 (2015)  
672 105–113. doi:10.1016/j.buildenv.2015.03.015.
- 673 [33] H. Karatas, T. Derbentli, Natural convection and radiation in rectangular cavities with one  
674 active vertical wall, *Int. J. Therm. Sci.* 123 (2018) 129–139.
- 675 [34] H. Karatas, T. Derbentli, Natural convection in rectangular cavities with one active vertical  
676 wall, *Int. J. Heat Mass Transf.* 105 (2017) 305–315.
- 677 [35] H. Karatas, T. Derbentli, Three-dimensional natural convection and radiation in a rectangular  
678 cavity with one active vertical wall, *Exp. Therm. Fluid Sci.* 88 (2017) 277–287.  
679 doi:10.1016/j.expthermflusci.2017.05.025.
- 680 [36] TL. Bergman, AS. Lavine, FP. Incropera, DP. DeWitt, *Fundamentals of Heat and Mass*  
681 *Transfer: Seventh Edition*, John Wiley & Sons, 2011.
- 682 [37] YF. Liu, ZY. Gao, XJ. Gao, *Heat transfer*, China Electric Power Press, 2015.
- 683 [38] G. Hetsroni, T.A. Kowalewski, B. Hu, A. Mosyak, Tracking of coherent thermal structures on  
684 a heated wall by means of infrared thermography, 30 (2001) 286–294.
- 685 [39] Y. Cheng, Z. Lin, Experimental study of airflow characteristics of stratum ventilation in a  
686 multi-occupant room with comparison to mixing ventilation and displacement ventilation,  
687 *Indoor Air.* 25 (2015) 662–671. doi:10.1111/ina.12188.
- 688 [40] Q. Ouyang, W. Dai, H. Li, Y.Z. Ā, Study on dynamic characteristics of natural and  
689 mechanical wind in built environment using spectral analysis, 41 (2006) 418–426..
- 690 [41] M. Lesieur. *Turbulence in fluids*, Springer, 2008.
- 691 [42] X. Zhang, J. Yu, G. Su, Z. Yao, P. Hao, F. He, Statistical analysis of turbulent thermal free  
692 convection over a small heat source in a large enclosed cavity, *Appl. Therm. Eng.* 93 (2016)  
693 446–455. doi:10.1016/j.applthermaleng.2015.10.011.