

Evaluating surface eddy properties in coupled climate simulations with 'eddy-present' and 'eddy-rich' ocean resolution

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1 **Evaluating surface eddy properties in coupled climate simulations with 'eddy-present' and 'eddy-**
2 **rich' ocean resolution**

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10 **Highlights**

11 • Eddy-rich (ER) has smaller and longer-lasting eddies than eddy-present (EP)

12 • EP captures 40% of eddies in observations even at high latitudes (ER captures 63%)

13 • Both model resolutions have a low eddy count in the EBUS and gyre interiors

14 • Eddy radii scale well with the minimum of the Rossby radius or the Rhines Scale

15 **Keywords** Mesoscale eddies; Eddy properties; High resolution coupled global model; Eddy dynamics

16 **Abstract**

17 As climate models move towards higher resolution, their ocean components are now able to explicitly re-
18 solve mesoscale eddies. High resolution for ocean models is roughly classified into eddy-present (EP, $\sim 1/4^\circ$)
19 and eddy-rich (ER, $\sim 1/12^\circ$) resolution. The cost-benefit of ER resolution over EP resolution remains debated.
20 To inform this discussion, we quantify and compare the surface properties of coherent mesoscale eddies in
21 high-resolution versions of the HadGEM3-GC3.1 coupled climate model, using an eddy tracking algorithm.
22 The modelled properties are compared to altimeter observations. Relative to EP, ER resolution simulates more
23 (+60%) and longer-lasting (+23%) eddies, in better agreement with observations. The representation of eddies
24 in Western Boundary Currents (WBC) and the Southern Ocean compares well with observations at both reso-
25 lutions. However a common deficiency in the models is the low eddy population in subtropical gyre interiors,
26 which reflects model biases at the Eastern Boundary Upwelling Systems and at the Indonesian outflow, where
27 most of these eddies are generated in observations. Despite a grid spacing larger than the Rossby radius of
28 deformation at high-latitudes, EP resolution does allow for eddy growth in these regions, although at a lower
29 rate than seen in observations and ER resolution. A key finding of our analysis is the large differences in eddy
30 size across the two resolutions and observations: the median speed-based radius increases from 14 km at ER
31 resolution to 32 km at EP resolution, compared with 48 km in observations. It is likely that observed radii
32 are biased high by the effective resolution of the gridded altimeter dataset due to post-processing. Our results
33 highlight the limitations of the altimeter products and the required caution when employed for understanding
34 eddy dynamics and developing eddy parameterizations.

35 1 Introduction

36 Mesoscale ocean eddies, generated from baroclinic and barotropic instabilities of the mean flow, are ubiquitous
37 in the world oceans [16]. Ocean eddies are important for a number of local processes such as air-sea exchanges
38 of momentum, freshwater and heat fluxes [60, 61, 33, 82] and the upwelling of nutrients, which promotes
39 biological activity [30, 8]. Mesoscale eddies have a major influence on the large-scale circulation, controlling its
40 mean state in the Antarctic Circumpolar Current [52], as well as its response to climate change [e.g. 2, 58, 34].

41 Over the last decade or so, many climate modelling groups have sought to increase the resolution of ocean
42 models [e.g. 55, 34, 67]. The primary aim has been to improve the representation of key mesoscale features
43 such as eddies, boundary currents and narrow sills (for dense overflows), and hence improve the mean-state and
44 variability of the coupled climate system [64, 57, 53]. It remains unclear whether the improved model fidelity
45 in higher resolution models is primarily a result of an improved mean state via these key frontal features, or a
46 consequence of the improved representation of the eddies themselves. The computational expense of a high-
47 resolution ocean component in a coupled climate model is high and the benefits of increased computational
48 cost need to be clearly identified.

49 In this context, the "high resolution" ocean component often refers to two types of resolutions: eddy-
50 present (EP, $\sim 1/4^\circ$) and eddy-rich (ER, $\sim 1/12^\circ$) [27]. Although not strictly defined, EP denotes resolutions
51 which permit some mesoscale eddies to be captured in the low and mid-latitudes, while ER refers to resolution
52 for which eddies are present at most latitudes (excluding the Arctic basin and the continental shelf around
53 Antarctica). The distribution of mesoscale features in a model mainly depends on the ratio of its horizontal
54 grid resolution, Δx , to the Rossby radius of deformation, R_d . Barotropic and baroclinic instability processes
55 are only expected to be properly resolved when the grid point spacing Δx is several times smaller than R_d ,
56 although a minimal criteria of 2 times smaller has sometimes been used [37].

57 Although coupled models with a high-resolution ocean component are increasingly available, many mod-
58 elling centres have not yet developed an operational version of their climate models with a high-resolution ocean
59 component. The upcoming Coupled Model Inter-comparison Project (CMIP6) will encompass models across
60 a range of resolutions, including EP and ER resolutions (CMIP6 HighResMIP) as well as eddy-parameterising
61 models [25, 36]. The cost-benefit balance of ER versus EP resolution is still being examined. While EP offers a
62 lower computational cost than ER resolution, it sits in the so-called 'grey-zone' where the benefits of removing
63 eddy parameterization and resolving some (but not all) mesoscale eddies and eddy fluxes are not obviously
64 superior to a coarser resolution ocean with full eddy parameterization [40]. Although the mesoscale field com-
65 prises more than just coherent eddies, evaluating the representation of coherent eddies at EP and ER resolutions
66 can inform the choice of resolution in future model development [40, 77].

67 Understanding the properties of eddies is also essential for their parameterization in coarse ocean models
68 [32]. For example, the eddy scale (estimated from either observations or models) often explicitly enters eddy
69 parameterization schemes through mixing length arguments e.g.[23, 6]. The size of coherent mesoscale eddies

70 is often used as an indicator of scale for the whole mesoscale field and is a fundamental measure employed in
71 numerous studies of eddy dynamics, notably to distinguish dynamical regimes [79, 78, 22, 75, 46].

72 While ocean models are not perfect tools to provide estimates of eddy properties, the robustness of the
73 spatial and temporal eddy scales from satellite altimetry has been questioned [13, 14, 17, 22]. Distortion of
74 the data can occur through the smoothing and interpolation required to generate a gridded product from raw
75 satellite measurements. Whilst high-resolution altimeters are currently being developed e.g. the future Surface
76 Water Ocean Topography (SWOT) mission, numerical simulations can allow us to evaluate eddy properties at
77 a much higher resolution than currently possible through observations [44, 80].

78 To date eddy properties have been studied in (coupled or ocean-only) high-resolution models at a regional
79 scale. Particular regions of interest include the Agulhas eddy pathways, important for heat transfer into the
80 South Atlantic [55], and the Californian Current System where eddies play a role in the transfer of heat and
81 nutrients from upwelling systems into the open ocean [48, 28]. Here we present a first global assessment of
82 mesoscale eddy properties (e.g. distribution, size, speed and lifetime) in two versions of the coupled model
83 HadGEM3-GC3.1 with EP and ER ocean resolution. Our study focuses on the field of coherent mesoscale
84 vortices, defined by closed sea surface height (SSH) contours, rather than the general mesoscale field compris-
85 ing filaments and unclosed structures. The characteristics (e.g. eddy kinetic energy, heat transport) of the two
86 fields likely differs [e.g. 74, 71]. We will address three central questions in this study: 1. As ocean resolution
87 in coupled models is increased, how does the representation of eddies and their properties change? 2. How
88 do modelled eddies and their properties compare to observations? 3. How do modelled eddies compare to
89 theoretical predictions?

90 This paper is organised as follows. Section 2 describes the eddy detection algorithm, and the model outputs
91 and observational datasets used. Section 3 presents results of global eddy counts and properties. Section 4
92 concludes and discusses the wider implications of the results.

93 **2 Method and Data**

94 **2.1 Eddy Detection Algorithm**

95 In this study, we use an eddy detection algorithm adapted from [54] (itself based on [14]). Eddies are identified
96 and tracked as closed coherent vortices detected through successive closed contours of SSH anomalies, subject
97 to various tests. The SSH field has a long-term 20 year mean removed. Large-scale SSH variability is removed
98 using a Gaussian filter with widths of $20^\circ \times 10^\circ$ (zonal \times meridional). The differences between this algorithm
99 and the original eddy detection algorithm of [14] are discussed in [54]. For example, this algorithm uses
100 interpolated SSH contours instead of raw SSH pixels, it includes a 'shape test' (to test how circular the closed
101 contour of SSH is) and a test for one local SSH minimum/maximum per eddy. Although the elongation of eddy
102 shape can play a role in the strength and extent of Western Boundary Currents (WBC), it is excluded from this
103 study [84]. Details of the scheme, criteria and tracking along with our adaptations of the filtering and detection

104 algorithm are further discussed in Appendix 1.

105 For both models and observations, the eddies are tracked globally using 20 years of daily SSH anomalies.
106 We only consider eddies with a minimum lifetime of 7 days. To minimize noise, the maps and probability
107 density functions (pdfs) of eddy statistics shown below only use eddies lasting longer than 1 month (unless
108 otherwise specified). Eddy properties considered in this study are as follows [14, 54]. The effective radius,
109 L_{eff} , is defined as the radius of a circle with the same area as the area within the outermost closed SSH contour
110 (satisfying all other criteria). The speed-based radius, L_{spd} , is taken as the radius of a circle similarly fitted to
111 the SSH contour with maximum averaged geostrophic velocity, U . By definition, L_{spd} is smaller than L_{eff} and
112 [14] found that typically $L_{spd} \simeq 0.7L_{eff}$. Eddy amplitude, A , is the absolute difference between the maximum
113 (for anti-cyclones) or minimum (cyclones) SSH within the eddy and the SSH value of the outermost closed
114 SSH contour (same contour as that used to define L_{eff}). The propagation velocity C_g^{eddy} is computed from
115 the daily displacements of the eddy center (defined as the center of a fitted circle to the smallest SSH contours,
116 i.e. 8 pixels). Here, we focus on the zonal component of C_g^{eddy} computed from the zonal displacements only.
117 Finally, a measure of eddy non-linearity is the ratio of the eddy rotational velocity to the eddy propagation
118 velocity, $r = U/C_g^{eddy}$. A value of r greater than 1 suggests that fluid parcels are trapped within an eddy [14].

119 There are numerous eddy detection algorithms available in the literature [86, 73, 49, 1, 15, 26, 29, 59].
120 They differ by the metric used for eddy identification (such as vorticity, Okubo-Weiss parameter or Lagrangian
121 particle tracking), filtering or the tracking technique (for example to include the merging and splitting of eddy
122 trajectories). Each method has its own advantages and limitations. The basis of this algorithm is physically-
123 based and has been heavily tested and used in literature [14]. In comparison to Lagrangian methods for example,
124 Eulerian tracking methods (such as closed SSH contours employed here) tend to over-estimate material con-
125 servation and transport, see [14, 1, 74]. However, a comparison of surface eddy properties is carried out here
126 instead of a quantification of eddy transport and energy. Crucially, in this study, the same eddy detection algo-
127 rithm is used on all datasets (model outputs and altimeter observations) to eliminate differences arising from
128 different detection algorithms.

129 **2.2 Coupled Model Configuration and outputs**

130 Outputs are analysed from the coupled high-resolution global climate model HadGEM3-GC3.1 [85]. This
131 model comprises a GA7.1/GL7.1 atmosphere/land configuration based on the MetUM and JULES [83], a GO6
132 ocean [70] based on NEMO [50] and GSI8 sea ice based on CICE [62]. Two resolutions of the ocean com-
133 ponent, both coupled to the same atmospheric component at resolution N216 (~ 60 km at mid-latitudes), are
134 compared: ORCA025 ($\sim 1/4^\circ$, hereafter EP_{sim}) and ORCA12 ($\sim 1/12^\circ$, hereafter ER_{sim}). The ocean com-
135 ponents do not employ any eddy parameterizations other than a small amount of isopycnal mixing to control
136 grid-scale noise. For further information about the model set-up, the reader is referred to [39, 85, 70].

137 The model simulations follow the CMIP6 HighResMIP protocol [36] with implementation described in
138 [63]. Model outputs (20 years of daily mean SSH) are obtained after a 20 year spin-up. Although the large-scale

continues to drift, it is likely that this has a negligible effect on eddy statistics, as changes in the background state are relatively small. In order to facilitate the comparison between versions of the coupled model, the eddy detection algorithm is also applied to 10 years of ER_{sim} SSH output re-gridded onto the EP_{sim} grid ($\sim 1/4^\circ$) (hereafter $ER_{sim}regrid$). The re-gridding was performed by bilinear interpolation, using an Earth System Modelling Framework (ESMF) [24], to generate conservative remapping of surface ocean variables (such as SSH) [39, 42].

2.3 Observational data

Observational SSH is taken from the gridded AVISO altimeter dataset (Archiving, Validation and Interpolating of Satellite Oceanographic Data, 2014; [21]). The Ssalto/Duacs altimeter products were produced and distributed by the Copernicus Marine and Environment Monitoring Service (CMEMS) (<http://www.marine.copernicus.eu>). The dataset provides daily SSH anomalies at $\sim 1/4^\circ$ resolution after the removal of a 20-yr mean. The gridded SSH field is generated through optimal interpolation from the delayed-time merging of multiple satellites. Note that, because we use an updated gridded altimeter product as well as a modified eddy tracking algorithm, our observed eddy statistics will differ from those published by [14].

Comparison of the raw daily SSH variances reveals differences before applying any filtering or eddy tracking, notably between observations and EP_{sim} . Although it captures the observed pattern correctly, EP_{sim} underestimates the magnitude of the observed variance, notably in WBC (not shown). ER_{sim} , however, compares reasonably well with observations on a global scale. Similar conclusions are obtained when comparing surface Eddy Kinetic Energy (not shown).

3 Results

We re-emphasize that the eddies detected in both observations and the model mostly consist of non-linear mesoscale coherent vortices in geostrophic balance. Most eddies in the ER_{sim} , EP_{sim} and observations have a small Rossby number R_o ($= \frac{U}{fL_{spd}}$): only 0.5, 0.06 and 0.09% of eddies in ER_{sim} , EP_{sim} and observations, respectively, have a Rossby number larger than 0.1 (Fig. A1, right). That is, none of the detected eddies, in the models or observations, are in submesoscale range (here we follow [56, 76] who define submesoscale as features with a Rossby number of order 1, among other criteria; this contrasts with other works which define submesoscales as smaller than 50 km [72]). Finally, as shown in Fig. A1 (left, note the logarithmic scale), most eddies have a non-linearity parameter r larger than 1.

3.1 Eddy Genesis and Lifetime

We start by comparing the rate and location of eddy genesis. Here, "eddy genesis" effectively refers the first time an eddy is identified. Although this is not the exact time when an eddy is born, this is a reasonable proxy. Fig. 1 shows maps of eddy genesis as the averaged frequency of first eddy detection in each 1° grid box per year.

171 Note that eddies require a minimum lifetime of 1 week to be identified by the detection algorithm. Differences
172 between models and observations are not sensitive to this choice – see the eddy genesis maps for eddies lasting
173 longer than 1 month in Fig. A2.

174 As expected eddies are not born homogeneously across the global ocean. Large genesis rates are found
175 in the vicinity of intense currents such as the Antarctic Circumpolar Current (ACC) and boundary currents.
176 Genesis rates are low in the open oceans, typically a factor of 4 smaller than in energetic regions. Model and
177 observations share broadly similar distributions of eddy genesis although the modelled rates are significantly
178 lower, notably in EP_{sim} . As a result, genesis rates in the gyre interiors of EP_{sim} approach zero. In addition
179 closer inspection reveals that genesis rates in EP_{sim} at Eastern Boundary Currents (EBCs) are very weak
180 compared to observations and ER_{sim} . This is particularly noticeable along the west coasts of Australia, Africa
181 and South America around 20-30°S. In contrast, ER_{sim} is able to capture these hot-spots of eddy genesis,
182 as well as generate as many eddies in the Southern Ocean as found in observations. This can be attributed
183 to improvements in the representation of ocean currents and outflows in ER_{sim} , partly through improved
184 topography, which provides a source of frontal shear for eddies to form [20]. For example improvements in
185 ER_{sim} are found in the Mediterranean outflow, EBCs, the ACC and the Drake passage, as well as in the East
186 Australian and Leeuwin currents around Australia [41]. However, ER_{sim} fails to capture the high genesis rates
187 of the North Atlantic and North Pacific sub-polar gyres as well as the long-lived (> 6 months) cyclonic eddies
188 from the Leeuwin Current and Tasman Outflow around Australia found in observations (see Fig. 2 below).

189 Table 1 shows the total number of eddies detected that last more than one week, as a crude measure of the
190 global eddy genesis. In all data sets the genesis rate are similar for cyclonic and anti-cyclonic eddies. However
191 consistent with Fig. 1, genesis rates are significantly lower in the models than in observations: ER_{sim} and
192 EP_{sim} generate only about 63% and 40% respectively of eddies found in observations. These biases in genesis
193 rate are reflected in the eddy counts for eddies with lifetimes longer than 4 weeks (even for eddies living more
194 than 16 weeks in EP_{sim}). For longer time-scales, other effects are playing a role (see below). These differences
195 in eddy genesis between the ER_{sim} and observations indicate that the ER resolution may still be too coarse
196 to generate mesoscale (coherent) eddies realistically. This may reflect that 1/12° (and 1/4°) resolution fails to
197 capture some smaller scale processes (e.g submesoscale activity, convection) that act as 'seeding' mechanisms
198 for the mesoscale activity through an inverse cascade of energy [65, 10, 56, 9].

199 Consistent with the genesis rates, the density of eddy tracks is larger in ER_{sim} and observations than in
200 EP_{sim} especially in eddy-energetic regions such as the Southern Ocean and WBCs (Fig. 2). For readability
201 only eddy trajectories lasting longer than 6 months are shown (the trajectories for all eddies lasting more than
202 2 months cover most of the ocean as shown in Fig. A3).

203 Eddies lasting longer than 6 months are concentrated in the subtropical gyres between 20° and 50° latitude.
204 They originate mainly from EBCs and to some extent from WBCs, notably from the Gulf Stream and North
205 Atlantic drift. Overall, the EP_{sim} significantly under-estimates the number of long lasting eddies although
206 anti-cyclonic eddies from the Agulhas current retroflexion ('Agulhas rings') are relatively well represented.

207 These trajectories form an important component of the meridional overturning circulation by controlling the
 208 quantity of heat and salt entering the North Atlantic [7]. However in other locations an artificially high number
 209 of eddy trajectories is found in the EP_{sim} , for example west of the Indonesian outflow (which may affect the
 210 Agulhas leakage [5]). A striking feature of observations is the absence of long lived eddies within and south
 211 of the ACC path (note that eddies are detected as far as 70°S , see Fig. A3). In contrast, in EP_{sim} and most
 212 notably in ER_{sim} , the ACC path is highlighted by the presence of numerous long-lived eddies.

213 These differences between the ER_{sim} , EP_{sim} and observations are reflected in the statistics of eddy life-
 214 time (Fig. 3). On average eddies in EP_{sim} and observations have shorter lifetimes than in ER_{sim} . The (normal-
 215 ized) probability density distributions of the eddy lifetimes are similar for EP_{sim} and observations but exhibit
 216 lower values than for ER_{sim} for lifetimes of 6 months and longer (Fig. 3, left).

217 Geographically, models and observations exhibit similar distributions of eddy lifetimes although, as ex-
 218 pected from Fig. 2 and 3, values in ER_{sim} are larger, with a global mean lifetime of 2 months compared to
 219 1.8 months in EP_{sim} and observations (Fig. 4). Eddy lifetimes are large in mid-latitudes ($20\text{-}50^\circ$) in all data
 220 sets, and large along the ACC pathway, notably in the Pacific sector in models. As highlighted by the zonal
 221 average (Fig. 3b), eddy lifetimes reach typically 2.2-2.4 months near $30\text{-}40^\circ\text{S}$ and fall to about 1.4-1.6 months
 222 at high latitudes and in the tropics. While models and observations show remarkable agreement in the Northern
 223 Hemisphere (Fig. 3b and Fig. 4), lifetimes in the models are consistently longer than in observations south of \sim
 224 40°S . Near 60°S , zonally averaged eddy lifetimes in EP_{sim} and ER_{sim} are (respectively) ~ 1.2 and ~ 1.4 times
 225 longer than in observations. At the highest latitudes, the presence of sea ice may partly explain the discrepancy
 226 as AVISO does not provide SSH data under sea while the models do [39]. However, the contrast between
 227 modelled and observed lifetimes is also clear in the core of the ACC which is ice-free all year long, suggesting
 228 other issues (see discussion below).

229 It is remarkable that, globally, the ER_{sim} simulates as many eddies with lifetimes >16 weeks as seen in
 230 observations (Table 1) despite a significantly lower genesis rate (by 37%). This implies that the "survival rate"
 231 of eddies is much larger in ER_{sim} than in observations (and EP_{sim}) (Table 1). The survival rate up to 4 weeks
 232 is quite similar across the three data sets. However it is 1.5 times larger in ER_{sim} than in observations at 16
 233 weeks and up to 3 times larger at 1 year. It is noteworthy that the survival rates of observations and EP_{sim} are
 234 very similar.

235 3.2 Propagation

236 Away from boundary currents and topography, eddies travel mainly in the zonal direction (Fig. 2). Theoretical
 237 predictions suggest that non-linear mesoscale eddies propagate westward with a velocity close to that of non-
 238 dispersive long baroclinic Rossby waves [19]. The theoretical Rossby wave phase speed in the long wave limit
 239 is given by $C_g^t = -\beta R_d^2$ where R_d is the Rossby radius of deformation. In the models, R_d is computed as
 240 $R_d = \int_{-H}^0 \frac{N}{|f|\pi} dz$ where $N(z) = \sqrt{-\frac{g}{\rho_0} \frac{d\rho}{dz}}$ (Brunt-Vaisala frequency) and f is the Coriolis parameter. For
 241 observations, we use the Rossby radius from Chelton et al. [11]. As found in previous studies [e.g. 12, 46], the

242 observed propagation speed of eddies, C_g^{eddy} , closely matches the Rossby wave speed, C_g^t , outside of the ACC
 243 (Fig. 5). Note that Fig. 5 shows C_g^t computed for the observed and modelled climatologies. At high-latitudes,
 244 the eddy propagation speed, C_g^{eddy} , approaches zero but increases towards the equator up to $\sim 10\text{-}12\text{ cm s}^{-1}$
 245 (westward). In the Southern Ocean however, eddies are carried eastward by the barotropic component of the
 246 ACC, resulting in a net eastward propagation speed of $\sim 1\text{ cm s}^{-1}$ [46]. Modelled zonal eddy propagation
 247 speed, C_g^{eddy} , in both ER_{sim} and EP_{sim} shows very good agreement with observations, including in the ACC
 248 (Fig. 5). This reflects the good climatology of the models (also evidenced by the similarity of the modelled and
 249 observed Rossby radius, not shown) as well as a good representation of the barotropic ACC in both models.

250 The co-location of global westward-propagating eddy trajectories longer than 6 months reveals the small
 251 equatorward drift of anticyclonic and poleward drift of cyclonic eddies (Fig. 6). Figure 6 flips the direction
 252 of propagation for NH and SH eddies so the positive latitudes are equatorward and the negative latitudes are
 253 poleward. For anti-cyclones (red), this meridional displacement increases from observations to EP_{sim} , and
 254 to ER_{sim} : the regression coefficients are 0.19 , 0.23 and $0.3^\circ/^\circ$ for observations, the EP_{sim} and ER_{sim} with
 255 R^2 values of 69%, 82% and 78% respectively. This means that anti-cyclonic eddies in the ER_{sim} are dis-
 256 placed by about 15° latitude for every 50° longitude travelled, whilst they are only displaced $\sim 10^\circ$ latitude in
 257 observations. Most of these long-lasting anti-cyclonic trajectories form part of the Agulhas rings. Compared
 258 to observations, a larger north-westward displacement of the Agulhas rings is also present in the stand-alone
 259 ocean component (Parallel Ocean program) of the Community Earth System Model, but this bias is reduced
 260 in the coupled simulations [55, 68]. This suggests that the representation of air-sea feedback over mesoscale
 261 eddies may influence their meridional migration. Although a similar number of eddies are plotted in Fig. 6,
 262 differences also partly reflect the longer eddy-lifetime found in the models (Fig. 2 and 3) with longer-lasting
 263 anti-cyclonic eddies found in the ER_{sim} compared to observations (and EP_{sim}). In ER_{sim} and EP_{sim} , the
 264 meridional drift is smaller for cyclones than anti-cyclones with regression coefficients of $0.16^\circ/^\circ$ and $0.15^\circ/^\circ$
 265 respectively (Fig. 6). However in observations, the meridional displacement is larger in cyclones ($0.23^\circ/^\circ$)
 266 than anti-cyclones ($0.19^\circ/^\circ$), and the displacement for each polarity is more symmetric than in the models.
 267 Many of these observed cyclonic trajectories are found in the Indian Ocean. These trajectories are absent from
 268 the models and may explain the asymmetric behavior found.

269 A simple measure of how "stationary" eddy are is given by the ratio $D/\overline{L_{eff}}$, where D is the absolute net
 270 zonal distance of propagation of an eddy and $\overline{L_{eff}}$ is its lifetime-averaged effective radius (Fig. 7). This ratio
 271 is simply a measure of the zonal displacement of eddies in units of "eddy radius". Maps of $D/\overline{L_{eff}}$ (Fig. 7)
 272 reveal that on average eddies are relatively stationary, moving by 3 or 4 times their radius. This is in contrast
 273 with the impression given by Figs. 2 and 6, which only include eddies longer than 6 months. Fig. 7 reflects that
 274 overwhelmingly eddies are short-lived, with life-times of about 2 months (see Fig. 3). $D/\overline{L_{eff}}$ varies mainly
 275 in the meridional direction, decreasing from 6-7 in the Tropics down to 1-2 at high latitudes, which primarily
 276 reflect variations of the propagation speed C_g^{eddy} . The 10 fold change in propagation speed between tropics
 277 and high-latitudes (Fig. 5) is somewhat reduced in $D/\overline{L_{eff}}$ due to the counter acting effect of changes in L_{eff}

278 (decreasing from the tropics to high latitudes, see below). Interestingly, the pattern of $D/\overline{L_{eff}}$ in ER_{sim} is
 279 less zonally symmetric than in EP_{sim} or observations, with enhanced values of $D/\overline{L_{eff}}$ in eddy-energetic
 280 regions such as the Agulhas Current Retroflexion, WBCs and along the ACC path. The latter feature notably is
 281 absent from observations, and reflects the smaller eddies detected in the ACC of ER_{sim} , which are not found in
 282 observations (see below). In the EP_{sim} , eddies are effectively more stationary than in ER_{sim} or observations
 283 almost everywhere. This bias may affect the ability of eddies at this resolution to transport and mix properties
 284 in the zonal direction.

285 3.3 Eddy Amplitude, Rotational Velocity and Radius

286 Distributions of eddy amplitude and rotational velocity are very similar between the three datasets although
 287 there is a hint that the distribution of amplitudes in ER_{sim} is narrower than in EP_{sim} and observations (Fig. 8a).
 288 Most eddies have amplitudes A between 1 and 5 cm with a median values of 2 cm.

289 Differences in rotational velocity U are more noticeable, although models and observations share similar
 290 distributions (Fig. 8b). The peak of the distribution is displaced toward larger values in ER_{sim} (6 cm s^{-1})
 291 compared with observations and EP_{sim} (4 cm s^{-1}). In the ER_{sim} , 19% of eddies have a velocity faster than
 292 14 cm s^{-1} (dotted line in Fig. 8b), whilst 14% do in the EP_{sim} and 13% in observations. In addition the
 293 fastest eddies in the EP_{sim} , at about 80 cm s^{-1} , are noticeably weaker than in the ER_{sim} and observations at
 294 $120\text{-}140 \text{ cm s}^{-1}$ (not shown).

295 The largest differences between the models and observations can be found when inspecting the radius of
 296 eddies (Fig. 9). Distributions are shown for both the speed-based L_{spd} and effective radii L_{eff} . The three
 297 distributions of eddy radius L_{spd} are very distinct, with median values of 48, 32 and 14 km for observations,
 298 EP_{sim} and ER_{sim} respectively. In the ER_{sim} , about a quarter (24%) of eddies have a radius L_{spd} equal to or
 299 smaller than 10 km while 90% of eddies have a radius L_{spd} equal to or smaller than 24 km (note that because
 300 of the convergence of the grid towards the poles, grid points can be significantly smaller than 10 km in ER_{sim} ;
 301 see Fig. 10). Instead 23% of eddies in the EP_{sim} and no eddies in observations have a radius L_{spd} equal to or
 302 smaller than 24 km. Conversely, both the ER_{sim} and EP_{sim} do not capture many eddies with a large L_{spd} :
 303 while in observations about 50% of eddies have a radius L_{spd} equal to or larger than 48 km, only about 6% in
 304 the EP_{sim} and 0.5% in the ER_{sim} reach such values.

305 Differences are less striking, but still significant, in terms of the effective radius L_{eff} (Fig. 9). EP_{sim}
 306 and ER_{sim} share similar distributions with median values of 52 km and 39 km, respectively. The observed
 307 distribution for L_{eff} is centred around 50 km but it is narrower than in EP_{sim} . It is interesting to observe
 308 that L_{eff} and L_{spd} are more similar in observations than in the models (Fig. 9). While L_{spd} is only slightly
 309 smaller than L_{eff} in observations (as in [14]), it is typically 2-2.5 times smaller than L_{eff} in the models. L_{spd}
 310 is likely to be much smaller than L_{eff} for a Gaussian-shaped eddy whereas the two measures should be nearly
 311 equal for a quadratic-shaped eddy [e.g. 14]. This may suggest that the profiles of observed eddies are closer
 312 to a quadratic shape while the profiles of modelled eddies better match a Gaussian shape. More likely, this

313 may reflect the large eddy radii found in observations. As the spatial scale grows in observations, closed SSH
 314 contours that satisfy the eddy algorithm criteria (e.g. no secondary extrema, shape test for circularity) are less
 315 likely to be found: L_{spd} is matched with L_{eff} (i.e. L_{spd} is reached at the edge of eddy).

316 3.4 Controls on Eddy Scales

317 Numerous studies have discussed processes that control the scale of ocean mesoscale eddies (e.g., [22, 78, 75,
 318 69, 47]). In this section, we discuss the eddy scales of coherent vortices in observations (based on an updated
 319 dataset compared to previous publications) and the EP_{sim} and ER_{sim} simulations in the light of these previous
 320 studies. The relationship between the size of eddies, the Rossby radius of deformation and the Rhines scale
 321 ($L_{Rhines} = \sqrt{\frac{U_{rms}}{2\beta}}$) is a recurring topic of investigation. A series of studies [75, 22, 79] have notably proposed
 322 that two regimes of ocean dynamics can be distinguished. They suggest that at low latitudes where L_{Rhines} is
 323 smaller than R_d , eddies scales with L_{Rhines} while at higher latitudes where L_{Rhines} is larger than R_d , eddies
 324 scales with R_d . The transition between the two regimes is found near 30°N/S (or $L_{Rhines} \simeq R_d \simeq 30$ km)
 325 equatorward of which baroclinic eddies can transfer their energy to Rossby waves [22, 79].

326 Starting with the models, it is interesting to note that the EP resolution allows eddy growth and propagation
 327 in high latitudes, as far as 60-70°N/S, where the EP grid scale is larger than the Rossby radius R_d . Following
 328 [37], Fig. 10 (top left) compares R_d with twice the grid scale Δx for the EP and ER resolutions. This criteria
 329 is inspired by linear stability analysis of baroclinic systems (e.g. the Eady and Charney problems; see [81]
 330 for a summary), which shows that maximum growth of linear waves is reached for wavelengths close to the
 331 Rossby radius of deformation. According to this simple criteria, eddies are expected to be found at nearly all
 332 latitudes in ER_{sim} but should be absent poleward of 30°N/S in EP_{sim} [38]. As evidenced by Figs. 1 and 2, this
 333 simple criteria does not apply in EP_{sim} . It is worth recollecting that although linear stability analysis predicts a
 334 maximum growth around the Rossby radius scale, it also predicts instability for a range of wavelength, including
 335 those larger than R_d . For example in the Eady problem, all wavelengths larger than $2.6 R_d$ are unstable while in
 336 Philip’s two-layer model, which includes a large wavelength cut-off due to the β -effect, unstable wavelengths
 337 are found between $2.2 R_d$ and $2\pi\sqrt{U_s/\beta}$ (where U_s is the mean vertical shear). We speculate that in regions
 338 where the grid scale is larger than the Rossby radius, instability and eddy growth remain possible but occur on
 339 scales significantly larger than the Rossby radius (or than the scale of the maximum theoretical growth rate).
 340 Indeed, most eddies in EP_{sim} (81%) are larger than R_d , unlike in ER_{sim} where only 20% are. This suggests
 341 that in the models (notably in ER_{sim}) the eddy scale is partly set by the grid scale or the smallest multiple of
 342 Δx that allows the development of instabilities. It should also be noted that C-grids (as used in NEMO) may
 343 develop a spurious baroclinic short-wave instability [4]. Such spurious mode may contribute to the smaller
 344 eddy scales found in ER_{sim} .

345 Further comparison reveals that the nominal and effective resolutions of these datasets, to be contrasted
 346 with the resolution of the underlying physics, also have a major influence on the estimated scales. To highlight
 347 this, the distribution of eddy scales for the ER resolution outputs are re-gridded to EP resolution (referred to as

348 $ER_{sim}regrid$) as shown in Fig. 9 (dotted lines). Through the remapping, the peak of the distributions for L_{spd}
349 in ER_{sim} increases from 12 km to 28 km. For L_{eff} , after remapping the distribution of ER_{sim} is shifted to
350 larger values by 12 km. Sensitivity tests with the high-pass filtering of the SSH field does not alter significantly
351 the eddy radius distributions (not shown). Not surprisingly, estimates of eddy scales are highly sensitive to the
352 resolution of the dataset. It is however striking that the distributions of eddy radii for $ER_{sim}regrid$ are nearly
353 identical to those for EP_{sim} . This reinforces the argument above that eddies grow on a scale set by the grid
354 scale. Despite having the same nominal grid resolution of $1/4^\circ$, observed eddy radii exhibit marked differences
355 with those of EP_{sim} and $ER_{sim}regrid$, notably for L_{spd} . If the re-mapping of ER_{sim} to $ER_{sim}regrid$ is any
356 guidance, this suggests that the effective resolution of the gridded observational dataset is larger than $1/4^\circ$ and
357 possibly closer to $1/2^\circ$.

358 A comparison of R_d , L_{Rhines} and L_{spd} is shown in Figs. 10 and 11. Equivalent plots for L_{eff} , which is
359 more noisy than L_{spd} , are shown in Fig. A4 and A5. Here we use $U_{rms} = \sqrt{EKE_g}$ to compute L_{Rhines}
360 where EKE_g is the surface geostrophic eddy kinetic energy (computed from 10 years of daily SSH anomalies
361 for the EP_{sim} and observations and from 5 years for the ER_{sim}). Note that L_{Rhines} is not defined in a standard
362 way in the literature. [22] uses the EKE associated with the barotropic flow. However, as the eddy velocity is
363 surface-intensified, our calculation of the Rhines scale is very similar to that of [22] (his Fig. 6). [79] define
364 L_{Rhines} as $2\pi\sqrt{\frac{2U_{rms}}{\beta}}$ and estimate U_{rms} as the root mean square of the eddy velocity from surface drifter
365 data. Since their U_{rms} and ours are similar (at least outside of the equatorial band, not shown), their estimate
366 of the Rhines scale for observations differs from ours (Fig. 10, top right) by a factor 4π . The Rhines scale is
367 similar for models and observations, ranging between approximately 30 and 60 km (Fig. 10)¹. Compared to
368 the Rossby radius R_d , L_{Rhines} exhibit a relatively flat, although noisy, meridional profile in all three datasets.
369 As the Rossby radius is also similar in models and observations, the ocean is separated in two regimes, with
370 $R_d \leq L_{Rhines}$ poleward of 30° N/S and $L_{Rhines} \leq R_d$ equatorward of 30° N/S.

371 The eddy radii vary quasi-linearly with latitude, increasing toward the equator (Fig. 10). Consistent with
372 Fig. 9, the zonally-averaged eddy radii are smallest in ER_{sim} and largest in observations. Again, eddy radii
373 in the regrided ER_{sim} is very similar to EP_{sim} (dark green line in Fig. 10). Note that the observed radii
374 L_{eff} (Fig. A4) compare well with Fig. 11 in [79] although our eddy radii are smaller. As the eddy detection
375 algorithm used in this study is essentially based on [15], this difference may be attributed to the fact that we use
376 a more recent altimeter product (with finer resolution).

377 Scatter plots of L_{spd} versus R_d or the minimum of R_d and L_{Rhines} are shown in Fig. 11 (see Fig. A5 for
378 L_{eff}). For observations and models, a good linear fit is found between L_{spd} (or L_{eff}) and R_d , although the
379 slope of the best fit between L_{spd} and R_d is slightly weaker in the ER_{sim} than in observations and the EP_{sim} :
380 for the ER_{sim} , EP_{sim} and observations slopes are 0.22, 0.35 and 0.35 with R^2 values of 90%, 80% and 82%

¹Note that there is no contradiction with the fact that L_{spd} differs substantially between models and observations as the Rhines scale and L_{spd} are not directly related. L_{Rhines} depends on the square root of the total geostrophic eddy kinetic energy while L_{spd} measures the distance between the eddy centre and the closed SSH contour with maximum averaged geostrophic velocity within an eddy.

381 respectively (Fig. 11, left column). For observations, this slope (0.35) is significantly smaller than the value of
 382 0.8 found in [22] in the North Atlantic while the fit found here appears much better than that seen in [22]. For
 383 both L_{spd} and L_{eff} , the relationship with R_d appears to break down (more scatter) for R_d larger than ~ 100 km
 384 (not shown). The scatter plots shown here are taken globally but a similar relationship is found for the North
 385 Atlantic only (see Fig. A6). Note however that [22] uses a different measure of the eddy size (based on the
 386 first zero-crossing of the spatial auto-correlation function of SSH anomalies) as well a older version of the SSH
 387 altimeter product.

388 Following [22], we test the relationship between the eddy radii and the minimum of R_d and L_{Rhines}
 389 (Fig. 11, A5 and A6, right column). The shade of colour indicates whether the minimum is reached with
 390 R_d (darker shade) or L_{Rhines} (lighter shade). In observations and EP_{sim} , the link between eddy radii and
 391 $\min(R_d, L_{Rhines})$ appears better than between eddy radii and R_d alone, as suggested in previous studies [e.g.
 392 22, 78]. Replacing R_d by $\min(R_d, L_{Rhines})$ clearly results in a more linear relationship to L_{spd} , as highlighted
 393 by the increased R^2 value, except from in ER_{sim} . However, this needs to be contrasted with the fact that the
 394 improvement of the fit (as measured by R^2) is often marginal and is sensitive to the choice of domain and of
 395 eddy radius definition (as shown for L_{eff} and for the North Atlantic in Figs. A5 and A6). Note that, as in [22],
 396 the slopes in EP_{sim} and observations are roughly double for $\min(R_d, L_{Rhines})$ relative to R_d .

397 4 Conclusions

398 Strengths and limitations of ocean simulations at ER and EP resolution in the representation of mesoscale
 399 eddies are explored. We focus on the surface properties of eddies using an eddy tracking algorithm on SSH
 400 anomalies. Modelled properties are compared to observed properties evaluated from the satellite altimeter
 401 AVISO product. An ocean model's ability to better-represent eddies in eddy-energetic regions, such as the
 402 WBCs, the Agulhas retroflexion and the Southern Ocean, has important implications for heat transport, global
 403 ocean stratification and eddy energy dissipation [58, 51, 88].

404 The key findings are summarized below:

- 405 • Amplitude, rotational speed and propagation speed of eddies are very similar across observations and
 406 models.
- 407 • ER and EP resolutions generate only $\sim 63\%$ and 40% respectively as many eddies as in observations.
 408 A leading factor for this discrepancy is the low count (or sometimes complete absence in EP) of eddy
 409 generation in the mid-ocean gyres and in Eastern Boundary Currents.
- 410 • Eddy lifetime are biased low in the EP_{sim} compared to observations but biased high in the ER_{sim} ,
 411 notably in the Southern ocean where the averaged eddy lifetime is about 30% larger than observed.
- 412 • Compared to EP_{sim} and observations, eddies are significantly smaller in ER_{sim} . This is true for both
 413 measures of eddy radius (speed-based and effective radius) although the differences are more striking for

414 the speed-based radius.

- 415 • Eddy radii scale closely with the Rossby radius of deformation, R_d , in all three datasets. As suggested
416 in previous studies, eddy sizes also relate well to the minimum of the R_d and the Rhines scale L_{Rhines} .
417 The improvement in the fit from R_d alone to $\min(R_d, L_{Rhines})$ is particularly notable in the ER_{sim} and
418 EP_{sim} .
- 419 • In contrast with suggestions from previous studies, EP_{sim} simulates a significant population of eddies
420 up to the high latitudes where the model grid-scale is larger than the Rossby radius of deformation, R_d .
421 These eddies likely grow on scales set by the smallest combination of grid-points that allows instability.

422 For the number of metrics explored in this study, it is difficult to objectively evaluate whether ER resolution
423 provides a significant improvement over EP resolution, in part due to concerns with the fact that observations
424 can provide a robust benchmark. Instead advantages of the ER_{sim} , compared to EP_{sim} , depend on the prop-
425 erties and region of interest. Benefits of the ER_{sim} include a similar number of eddies in the Southern Ocean,
426 and globally a similar number of eddies living longer than 16 weeks, compared to observations. ER_{sim} eddies
427 are less stationary and smaller eddies are able to develop, compared to the EP_{sim} . The genesis rate and size
428 of the eddy populations are clear examples where the ER_{sim} improves upon EP_{sim} . This is likely the result
429 of a better representation of the mean state in the ER_{sim} in eddy-energetic regions such as boundary currents
430 and the ACC. Eddies generated in Eastern Boundary Currents are important for transferring heat and nutrients
431 into the nutrient-poor open ocean [28, 31]. In that regard, the ER_{sim} clearly outperform the EP_{sim} where the
432 basin interior are relatively empty of eddies.

433 In other aspects, outcomes of the model-observation comparison are more ambiguous. Our results suggest
434 that the ER_{sim} over-estimates the survival rate of eddies. The dissipation of mesoscale eddies in the ocean
435 remains an open question with a number of competing ideas being explored e.g. enhanced friction over rough
436 bottom topography, the emission of internal waves, coupling to the atmosphere, the role of symmetric instability
437 in the open ocean or interaction with WBCs [88, 18, 35, 87]. It is not expected that such processes are captured
438 in ER (nor EP) resolution models. Our analysis suggests that as resolution increases, allowing more vigorous
439 eddies and a lower viscosity (for numerical stability), the absence of dissipation mechanisms may become
440 problematic and introduce biases in the lifetime of the modelled eddies. However, we cannot rule out that eddy
441 lifetime estimates are biased low in observations due to post-processing and smoothing of the SSH data that
442 would limit the ability to track eddies.

443 The differences in eddy size are a particularly striking outcome of our analysis. Our results suggest that
444 the eddy size is overestimated in observations by a factor 2 and possibly up to 4 depending on the considered
445 measure. The nominal resolution of the dataset is a key factor here and, consistent with previous studies, our
446 analysis suggests that the effective resolution of the AVISO gridded dataset is coarser than $1/4^\circ$ [13, 14, 73, 3].
447 Instead, the effective resolution in the ER_{sim} is much higher than in observations but the subsequent impact of
448 the smaller eddies found in the ER_{sim} is unclear. Whether the total energy or heat contained within a greater

449 number of smaller eddies in the ER_{sim} is similar to the fewer, larger eddies found in the EP_{sim} remains to be
450 determined. Further studies are needed to explore the role of the tracked eddies in air-sea and surface-subsurface
451 coupling within the climate system. An overestimation of eddy scales in observations could have implications
452 for eddy parameterization and interpretation of ocean dynamics. Mixing length arguments underlying many
453 eddy parameterizations use the eddy scale as proxy for the mixing length [45]. Direct comparison of properties
454 (e.g. wavenumber spectrum, see [66]) along satellite tracks should help clarify to which extent differences
455 between model and observations are robust or due to the post-processing necessary to generate the AVISO
456 gridded product.

457 Finally, it must be noted that our model represents one set of parameter choices, for example the sensi-
458 tivity to viscosity has not been tested, and only surface eddy properties are evaluated. Further studies should
459 explore the 3-dimensional structure of eddies, the influence of eddies of air-sea exchanges and energy spectra
460 to compare the redistribution of kinetic energy at larger scales for each resolution [43]. Limitations of the eddy
461 tracking algorithm should not be underestimated [14]. It is likely that some of our results (e.g. eddy counts)
462 are dependent on our choice of eddy detection algorithm. However we have attempted to minimize its impact
463 by applying the same algorithm to models and observations and focus our analysis on differences/similarities
464 rather than the absolute values. This work lays the foundation for future studies at different resolutions and
465 using different models as more high resolution data become available in which submesoscales start to be re-
466 solved. Observational SSH global datasets are likely to improve as satellite altimetry coverage is enhanced with
467 the future launch of the SWOT altimeter.

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472 **5 Appendix**

473 **5.1 Algorithm details**

474 Further discussion of the eddy identification and tracking scheme, criteria and adaptations are listed below.
475 SSH contours are computed from 100 cm to -100 cm with an interval of 0.3 cm. Starting from a SSH minimum
476 (cyclone) or maximum (anti-cyclone), the algorithm identifies successive closed contours. There is no set
477 minimum or maximum eddy radius; instead an eddy's size is limited by its pixel range. In order for an eddy
478 to be successfully identified each closed contour of SSH needs to lie within a specific pixel range between 8
479 and 10,000. Therefore when increasing the grid resolution the same minimum pixel number of 8 allow smaller
480 eddies to be detected compared to a coarser resolution.

481 Adaptations from the original eddy tracking algorithm [54] include:

- 482 • The identification and tracking components of the algorithm were split so global identification at each
483 daily timestep is run in parallel to increase computational efficiency. For a chosen region and time period,
484 eddies are then able to be tracked from the already identified eddy centres. All eddy tracks (and their
485 associated properties such as radius, rotational velocity and amplitude) are stored and for eddies left
486 'active' (not masked), their tracks are able to be resumed for future tracking.
- 487 • The regular grid is adapted for use with the irregular NEMO ocean grid. A remaining limitation to our
488 method is the ability to wrap tracks across the irregular NEMO grid divide at approximately 73°E. This
489 slight jump in tracks is assumed to not have a large consequence on global statistics and there is no
490 obvious increase in eddy birth and death frequency either side of this divide. This can be observed in
491 Figs. 1 and 7.
- 492 • Improvements were made in the unrealistic 'jumping' of eddy tracks by changing the search ellipse used
493 to find the following identified eddy contour in its track. This was based on tracking improvements pub-
494 lished online associated with a collaboration with AVISO [A. Delepouille et al. OSTST 2017, Mesoscale
495 Eddies in Altimeter Observations of SSH web site at OSU, [http : //wombat.coas.oregonstate.edu/eddies/](http://wombat.coas.oregonstate.edu/eddies/),
496 accessed 08.11.18.]

497 A link to the AVISO handbook can be found here with details of the tracking method: [https : //www.aviso.
498 altimetry.fr/fileadmin/documents/data/tools/hdbk_eddytrajectory_2.0exp.pdf](https://www.aviso.altimetry.fr/fileadmin/documents/data/tools/hdbk_eddytrajectory_2.0exp.pdf) accessed 20.03.19

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	Type	>1 wk	>4 wks	> 16wks	> 26wks	> 40wks	> 52wks	> 78wks
EP	A	143,944	29,721	2,099	495	96	41	13
	C	135,892	24,943	1,744	378	58	13	1
			<i>19.5%</i>	<i>1.4%</i>	<i>0.31%</i>	<i>0.06%</i>	<i>0.02%</i>	<i>0.005%</i>
ER	A	202,639	45,595	4,412	1,333	386	190	82
	C	205,633	41,642	4,003	1,240	346	155	33
			<i>21.4%</i>	<i>2.1%</i>	<i>0.63%</i>	<i>0.18%</i>	<i>0.08%</i>	<i>0.03%</i>
Obs	A	355,221	73,683	5,021	1,276	306	115	32
	C	334,599	64,064	3,874	933	206	70	11
			<i>20.0%</i>	<i>1.3%</i>	<i>0.3%</i>	<i>0.07%</i>	<i>0.03%</i>	<i>0.006%</i>

Table 1: Number of eddies detected with lifetimes longer than 1, 4, 16, 26, 40, 52 and 78 weeks for the eddy-permitting simulation EP_{sim} , the eddy-resolving simulation ER_{sim} , and the AVISO gridded satellite altimetry product (Obs). The counts are scaled to 10 years and separated for cyclonic (C) and anti-cyclonic (A) eddies. For each data set, the third line (in italic) indicates the survival rate, i.e. the ratio (expressed in %) between the total number of eddies with a given lifetime and the total number of eddies with lifetime longer than 1 week.

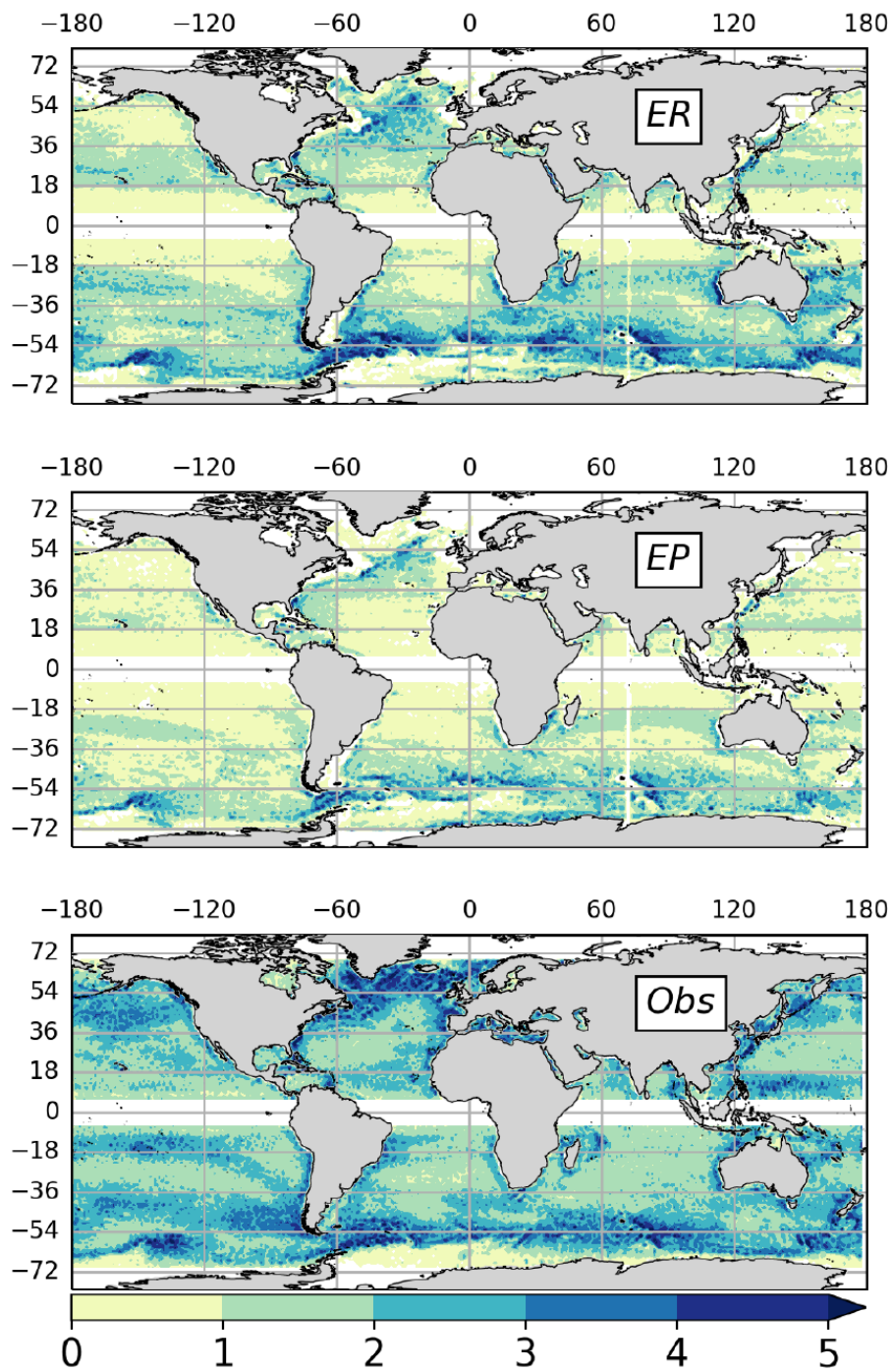


Figure 1: Eddy genesis (number of eddies per year) for eddies lasting longer than 1 week (binned to $1^\circ \times 1^\circ$ boxes).

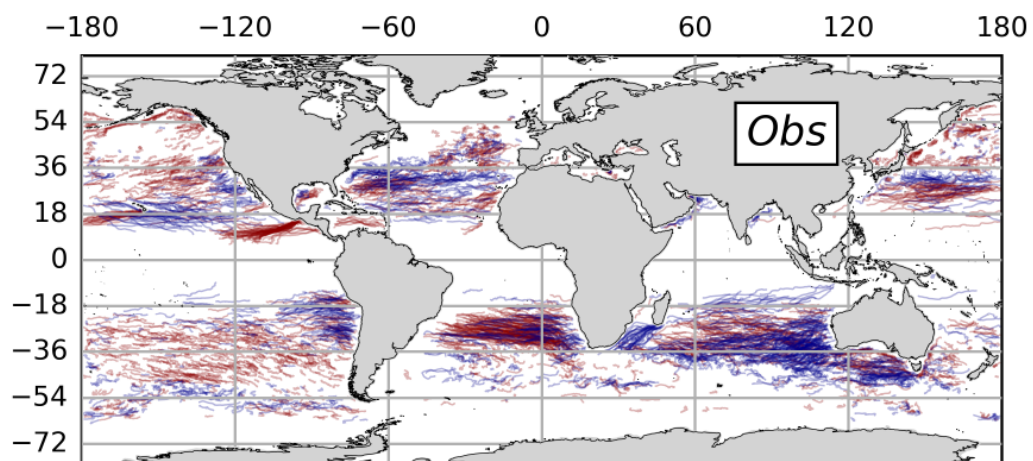
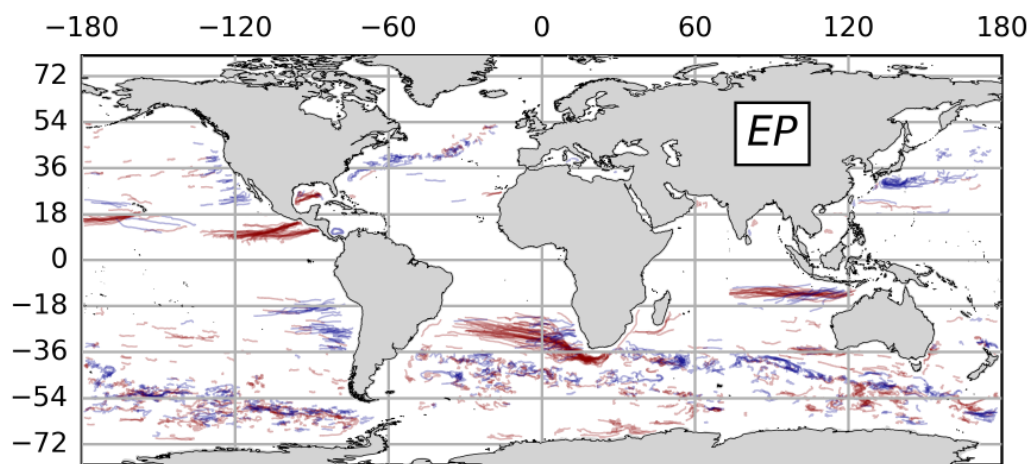
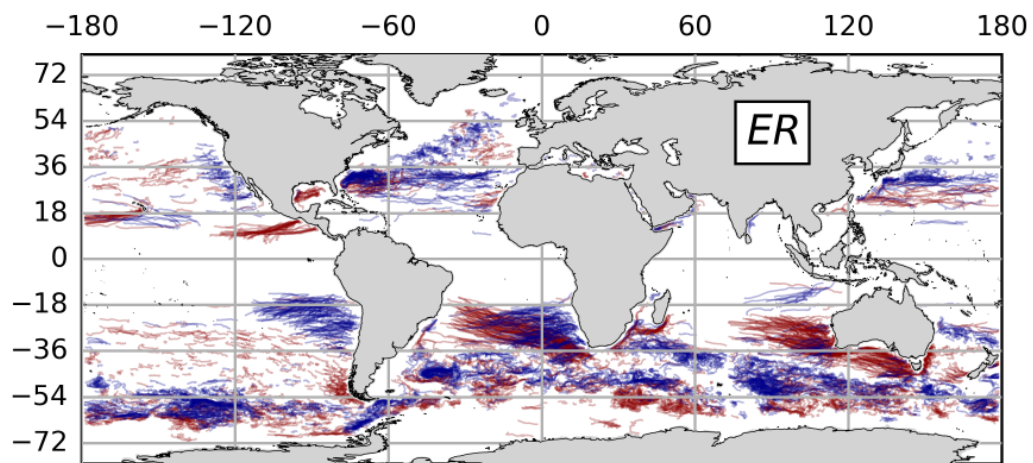


Figure 2: Eddy trajectories lasting longer than 6 months over 20 years. Anti-cyclonic (cyclonic) eddies are shown in red (blue).

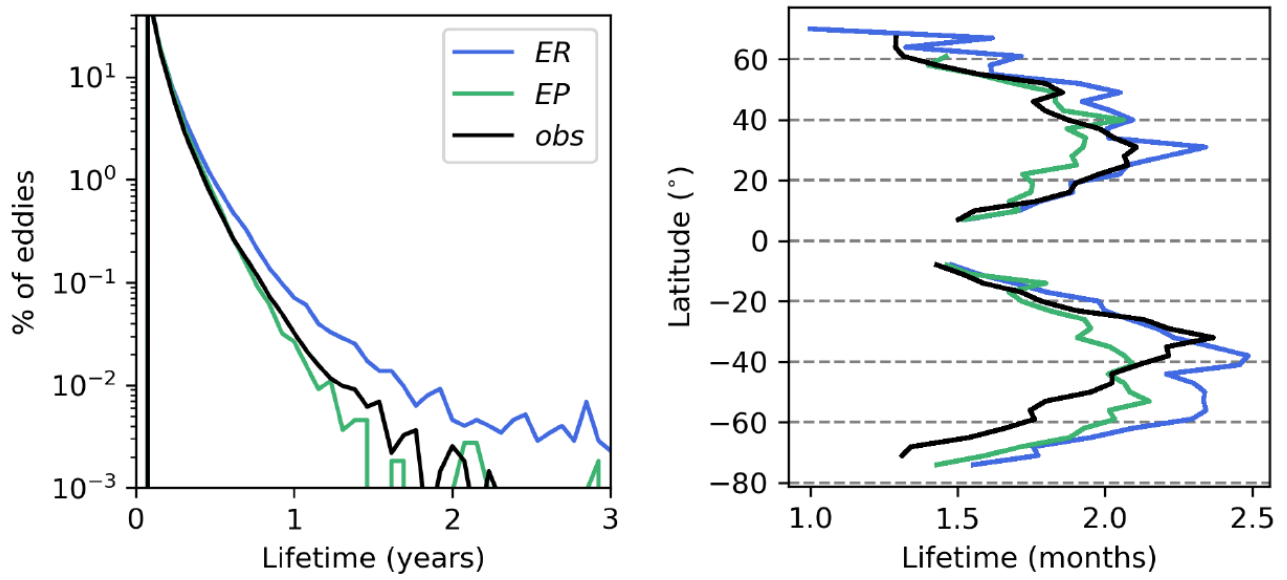


Figure 3: Probability density function of eddy lifetime (left) and zonal average of eddy lifetime (right). Both plots use eddies with lifetimes longer than 1 month.

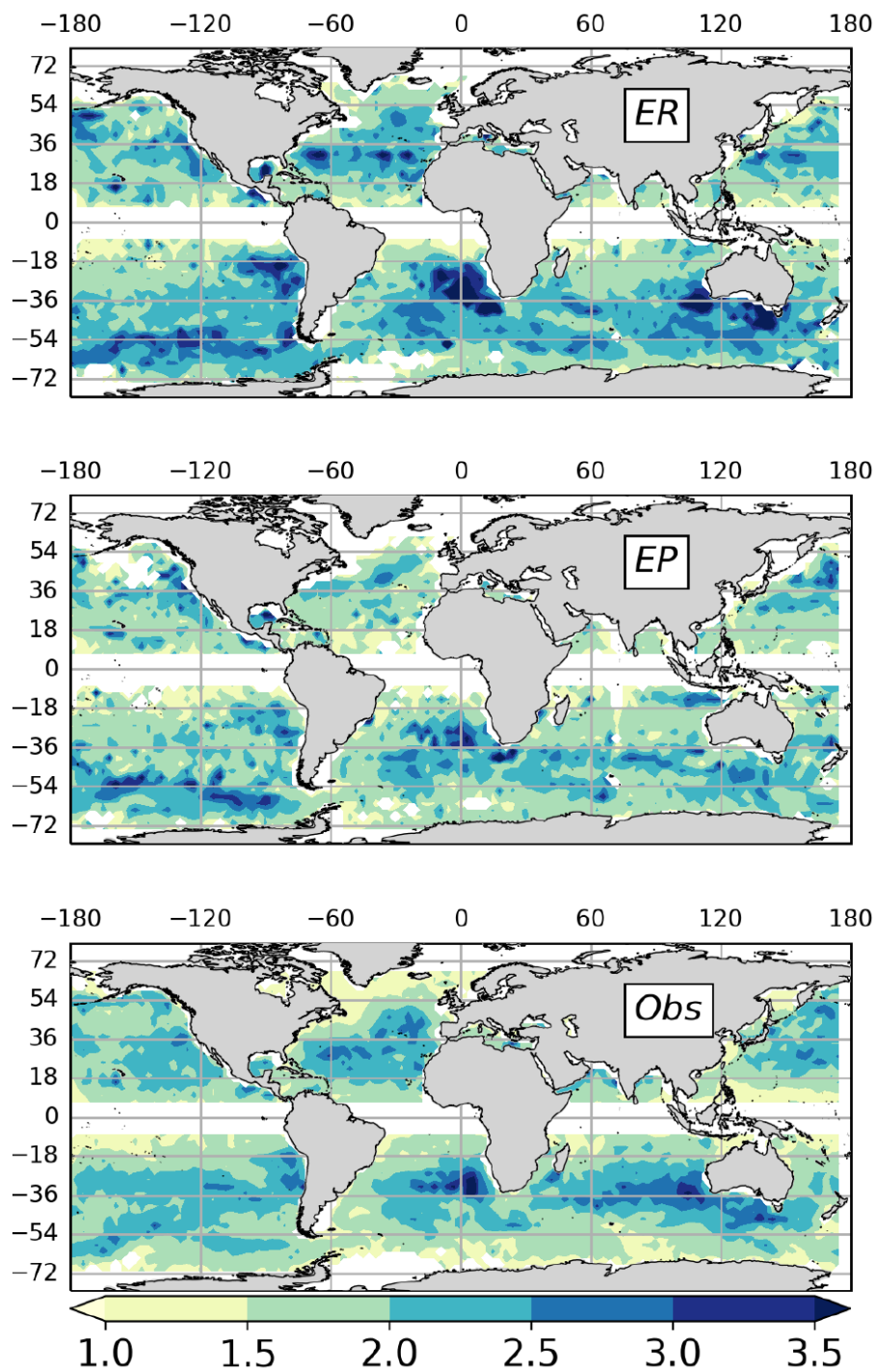


Figure 4: Eddy lifetimes (in months) mapped to genesis location and binned to $3^\circ \times 3^\circ$ grid boxes. All plots use eddies with lifetimes longer than 1 month.

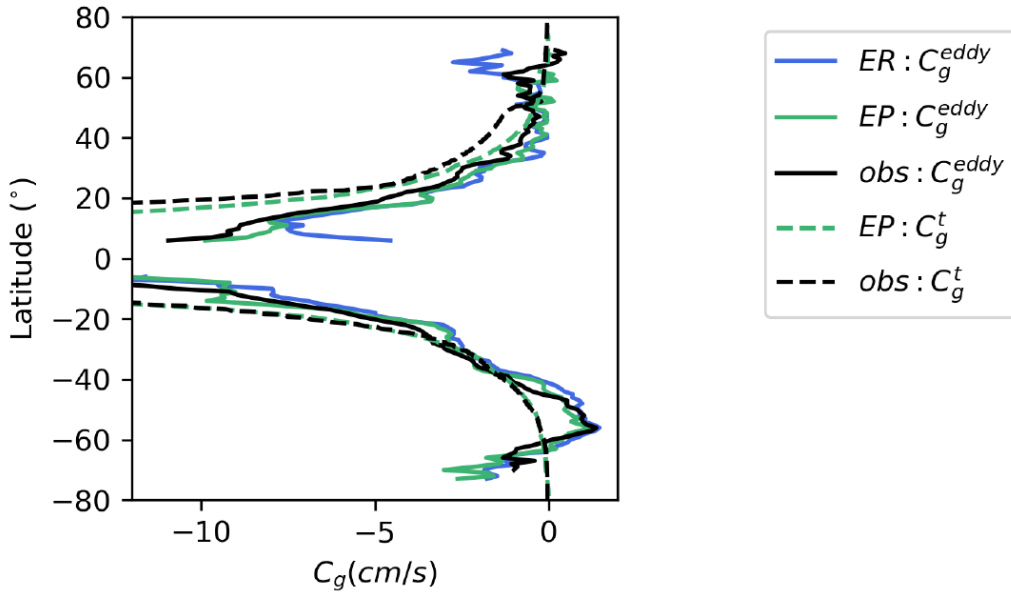


Figure 5: Zonal average of zonal propagation velocity (cm s^{-1}) from tracked eddies, C_g^{eddy} . Dotted lines are the theoretical long-wave baroclinic Rossby Wave speed C_g^t for observations (black) and $EP_{\text{sim}}/ER_{\text{sim}}$ (green).

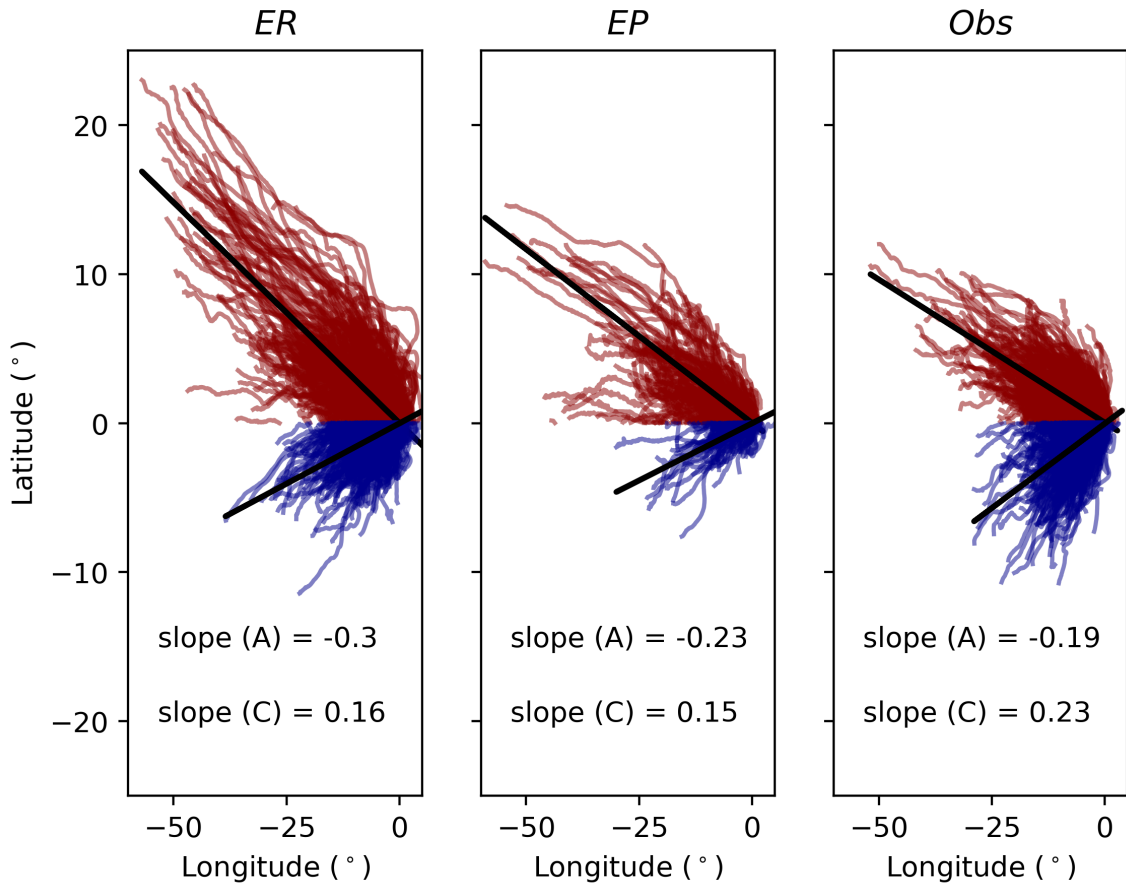


Figure 6: Co-located trajectories of westward-propagating eddies lasting longer than 6 months for ER, EP and observations. Anti-cyclonic eddies (A) are plotted in red, cyclonic eddies (C) are in blue and the regression coefficients for each are given on each subplot.

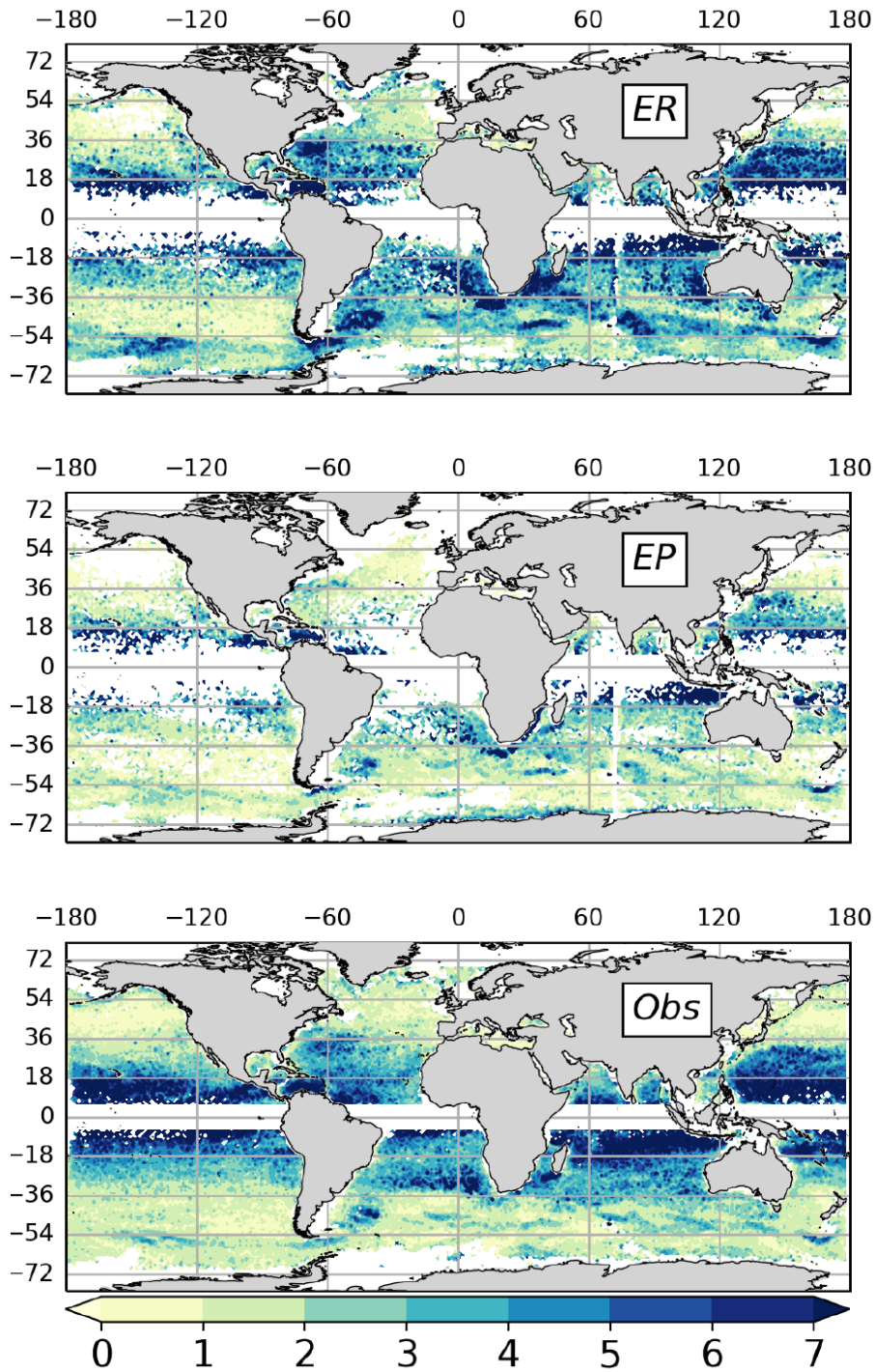


Figure 7: 20-year average of the ratio $D/\overline{L_{eff}}$ where D is net zonal distance covered by an eddy and $\overline{L_{eff}}$ its lifetime-averaged effective radius. The ratios are mapped to genesis locations and binned to $1^\circ \times 1^\circ$ boxes.

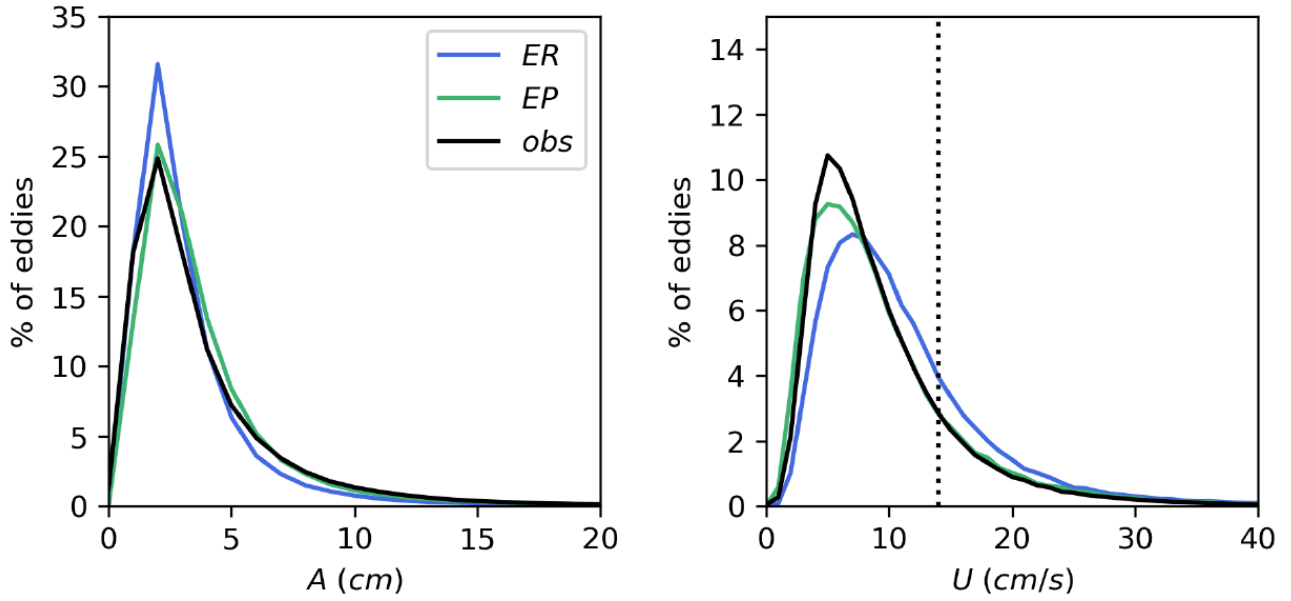


Figure 8: Probability density functions of the lifetime-averaged amplitude A (left) and rotational velocity U (right) of eddies longer than 1 month (with 1 cm and 1 cm s⁻¹ bins). The black dotted line is plotted at 14 cm s⁻¹.

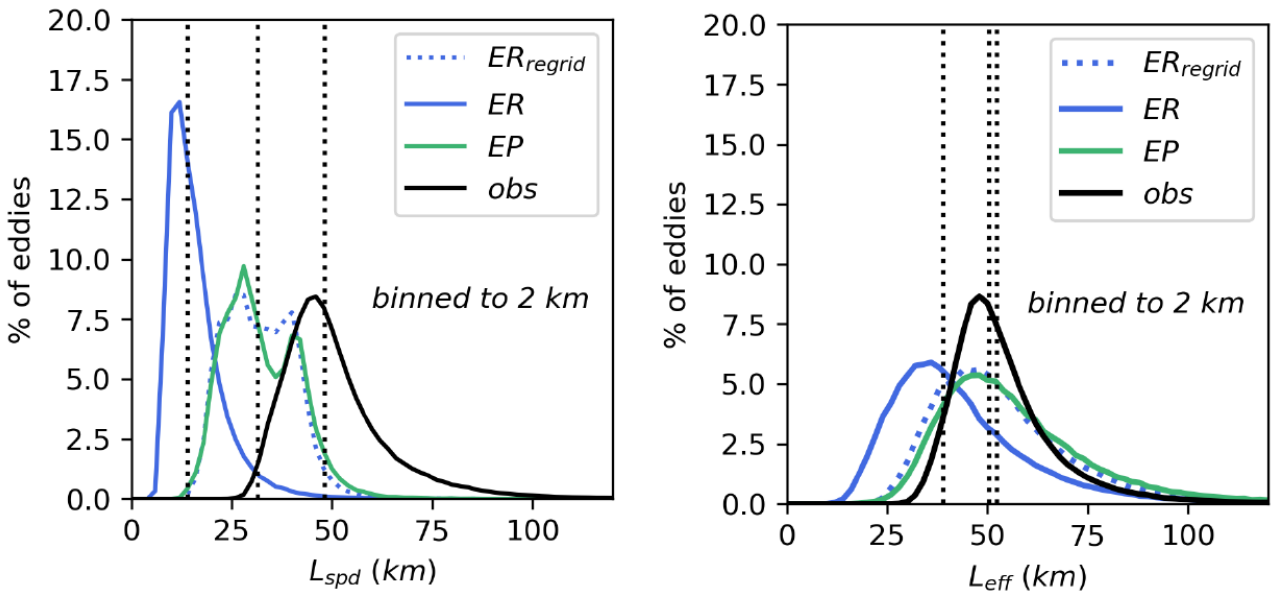


Figure 9: Probability density functions (pdfs) of the lifetime-averaged L_{spd} and L_{eff} : a normalized pdf on a linear scale with 2 km bins. The black dotted lines are plotted at the medians for each resolution: the median values for L_{spd}/L_{eff} are 48 km/50 km, 32 km/52 km and 14 km/39 km for observations, EP_{sim} and ER_{sim} , respectively.

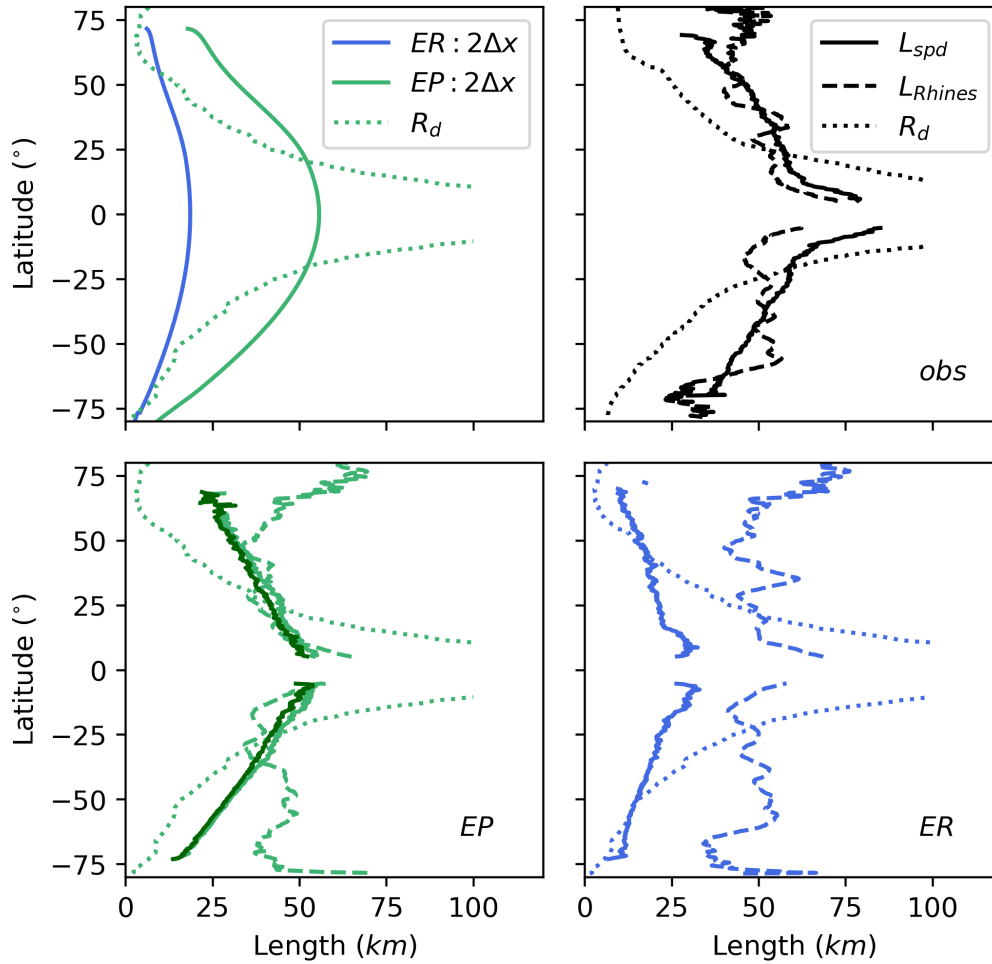


Figure 10: (top left) Zonal average of the observed Rossby radius of deformation R_d and $2\Delta x$ for EP and ER. (top right and lower subplots) Zonal average of L_{spd} (solid lines), the Rossby radius of deformation (R_d , dotted line) and the Rhines Scale (L_{Rhines} , dashed line) for observations (black), EP_{sim} (green) and ER_{sim} (blue). The zonal average of L_{spd} for ER_{regrid} is plotted in dark green.

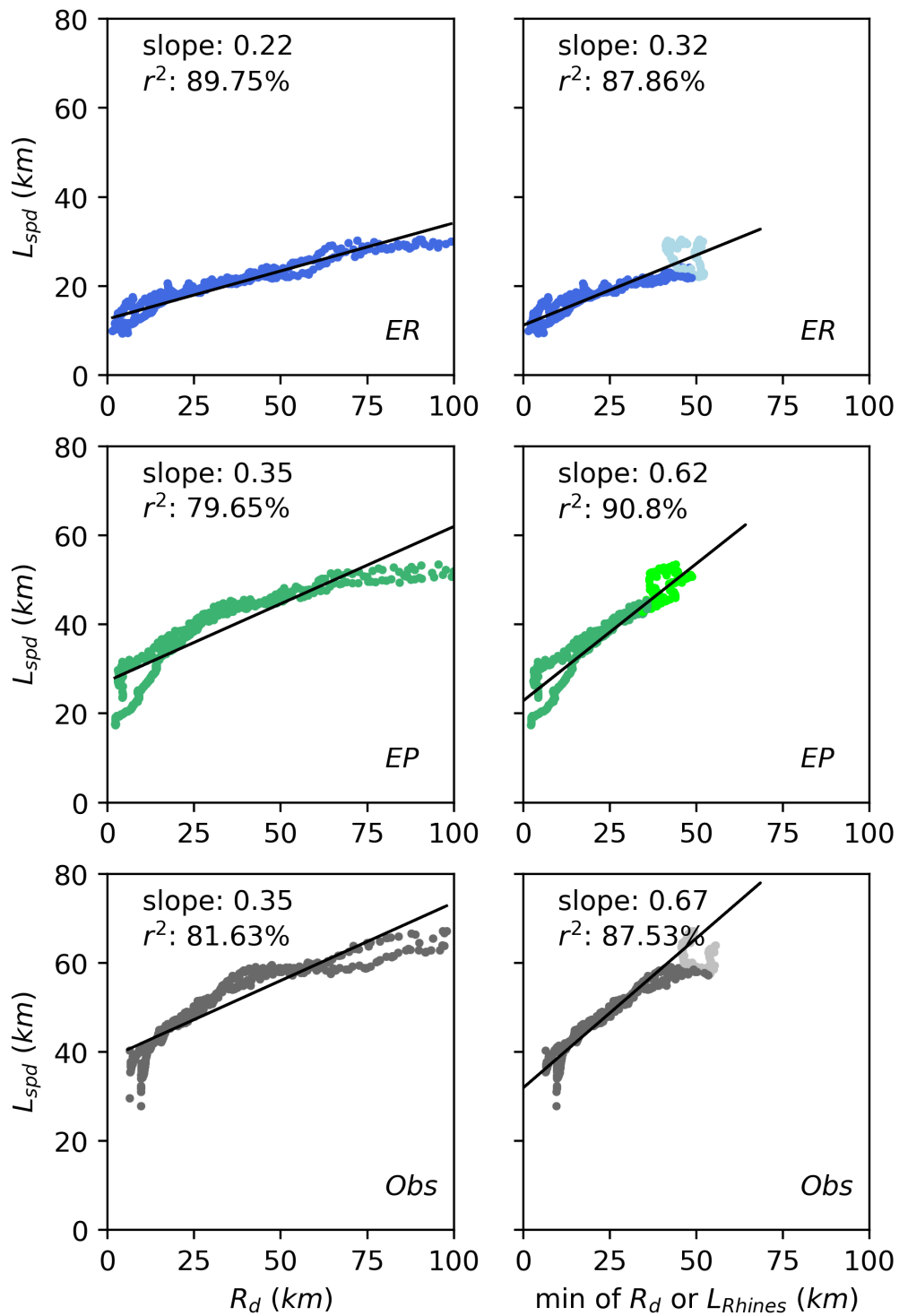


Figure 11: L_{spd} compared to R_d (left) and to the minimum of R_d and L_{Rhines} (right). The data is global after zonally averaging. The linear regression line is plotted in black. In the right panels, the shade of colour indicates whether the minimum is reached with R_d (darker shade) or L_{Rhines} (lighter shade). EP is plotted in blue, EP in green and observations are plotted in grey.

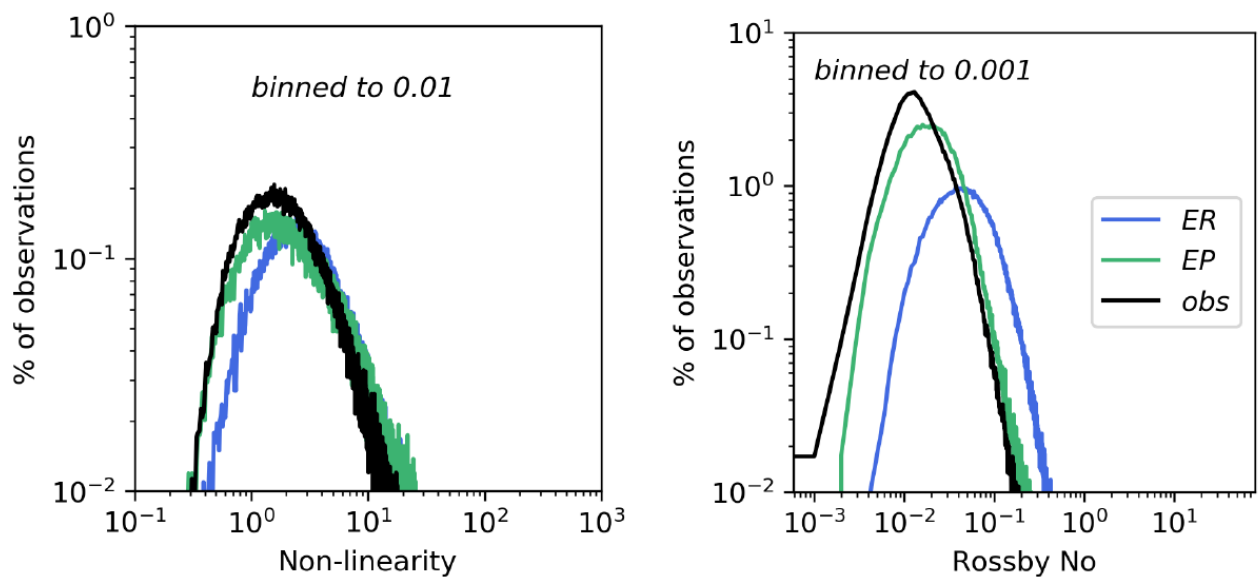


Figure A1: Probability density functions of (left) the non-linearity parameter r and (right) the Rossby number R_o .

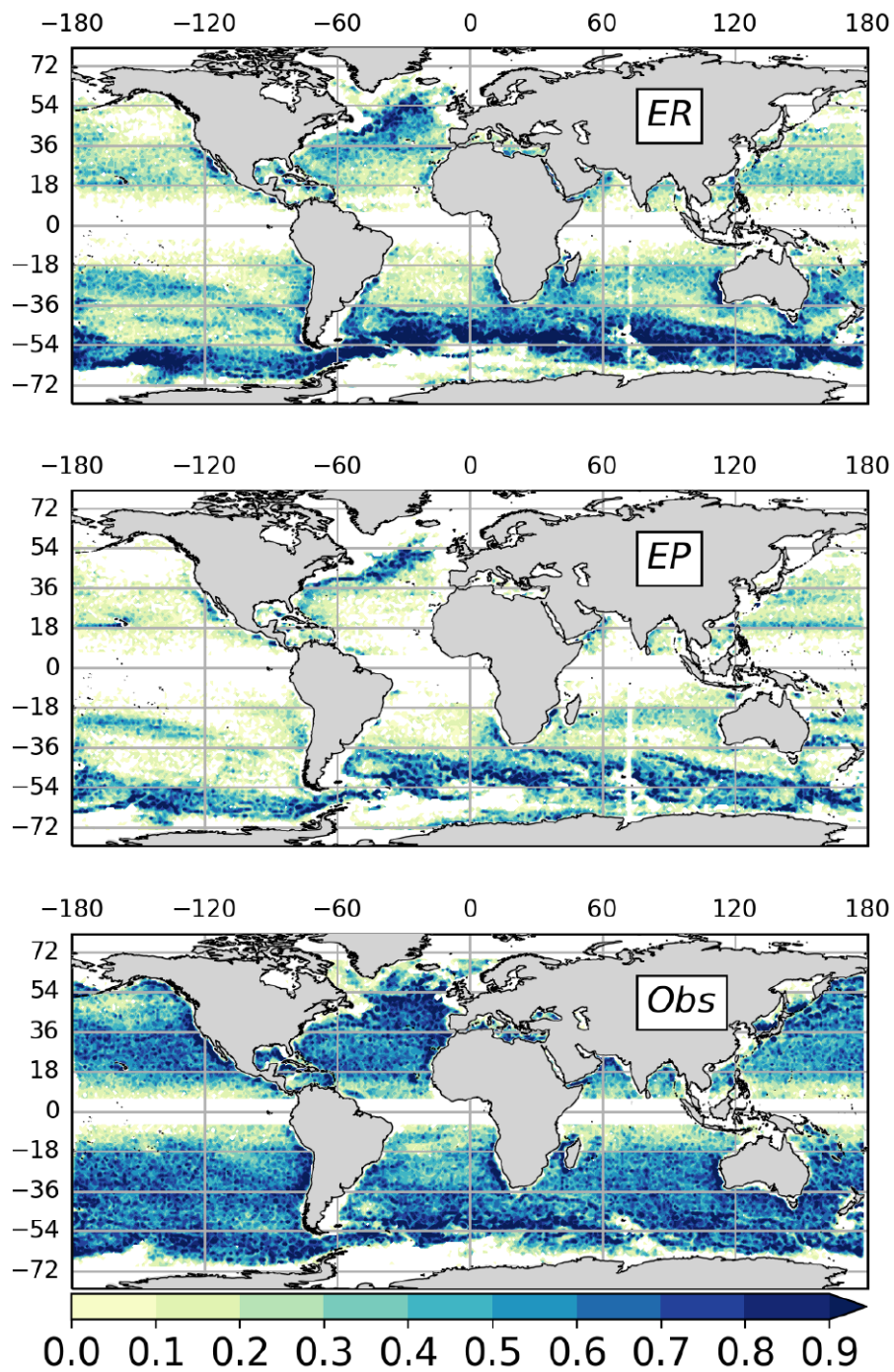


Figure A2: Eddy genesis (number of eddies per year) for eddies lasting longer than 1 month (binned to $1^\circ \times 1^\circ$ grid boxes).

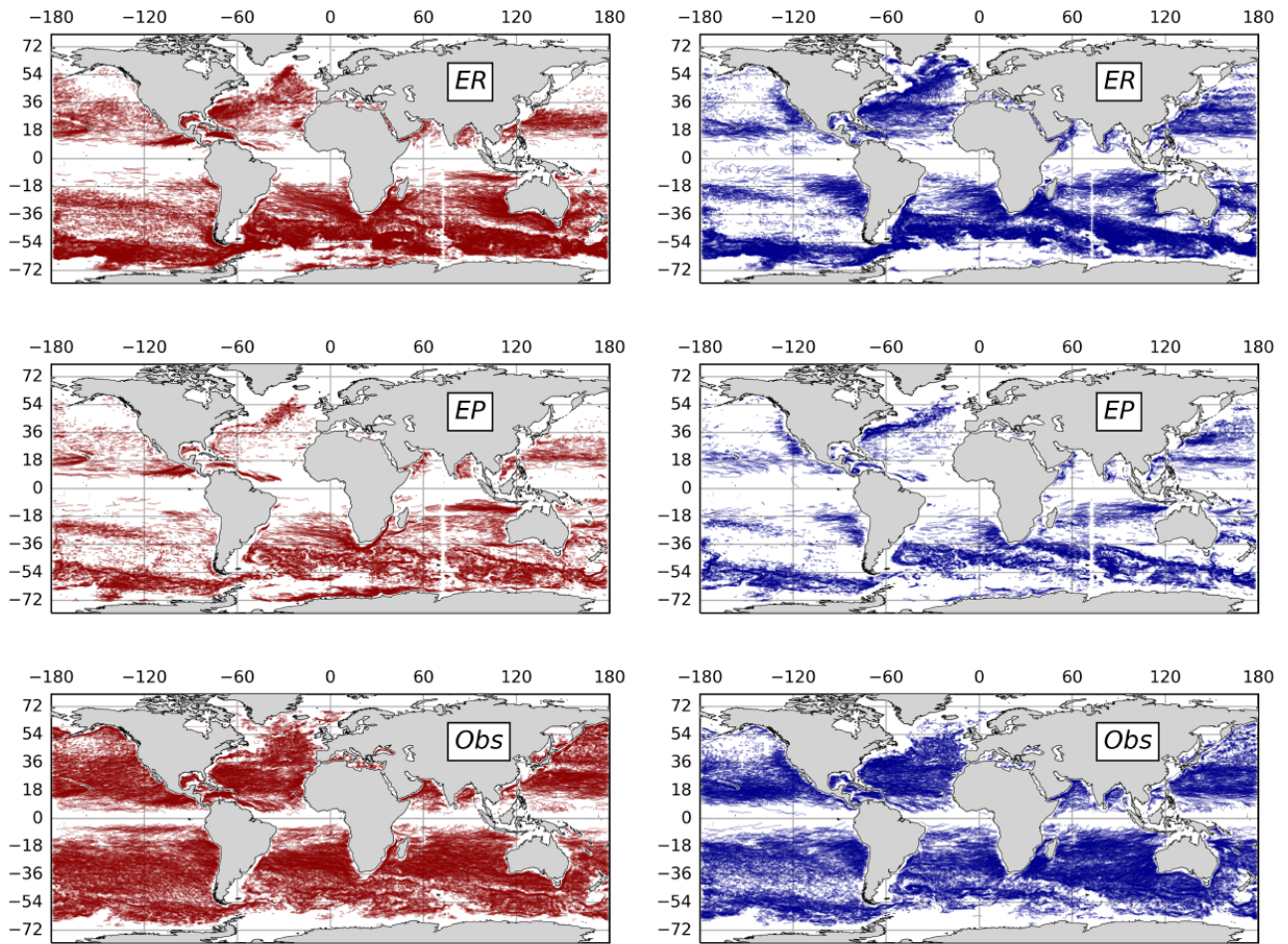


Figure A3: Eddy trajectories lasting longer than 2 months over 20 years. Anti-cyclonic eddies (left) are shown in red and cyclonic eddies (right) are in blue.

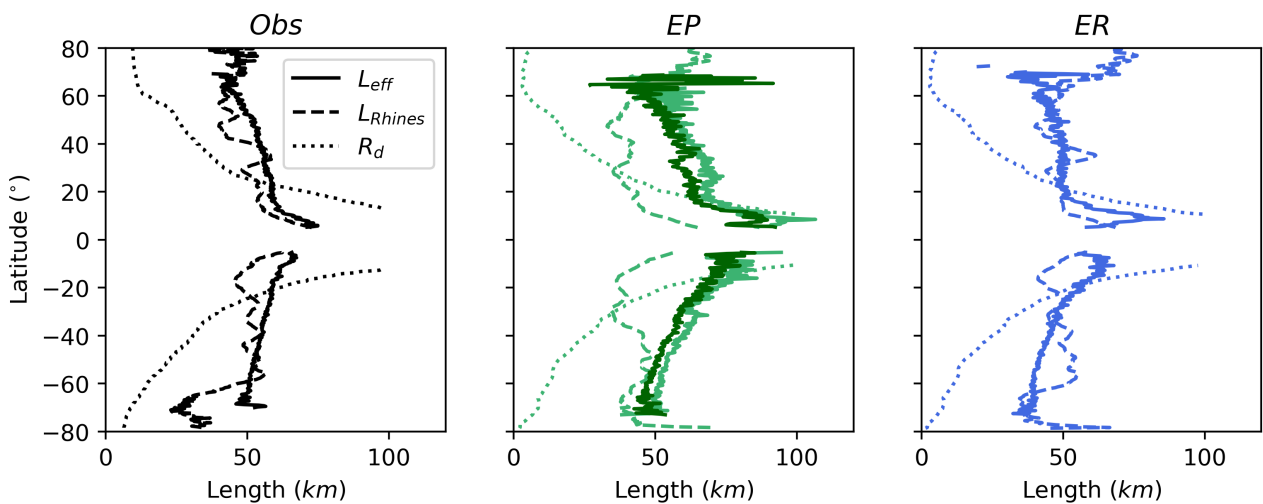


Figure A4: A repeat of the zonal average shown in Fig. 10 for L_{eff} (solid lines) against the Rossby radius of deformation (R_d , dotted line) and the Rhines scale (L_{Rhines} , dashed line). Observations are plotted in black, EP is in green and ER in blue. The zonal average of L_{spd} for ER_{regrid} is plotted in dark green.

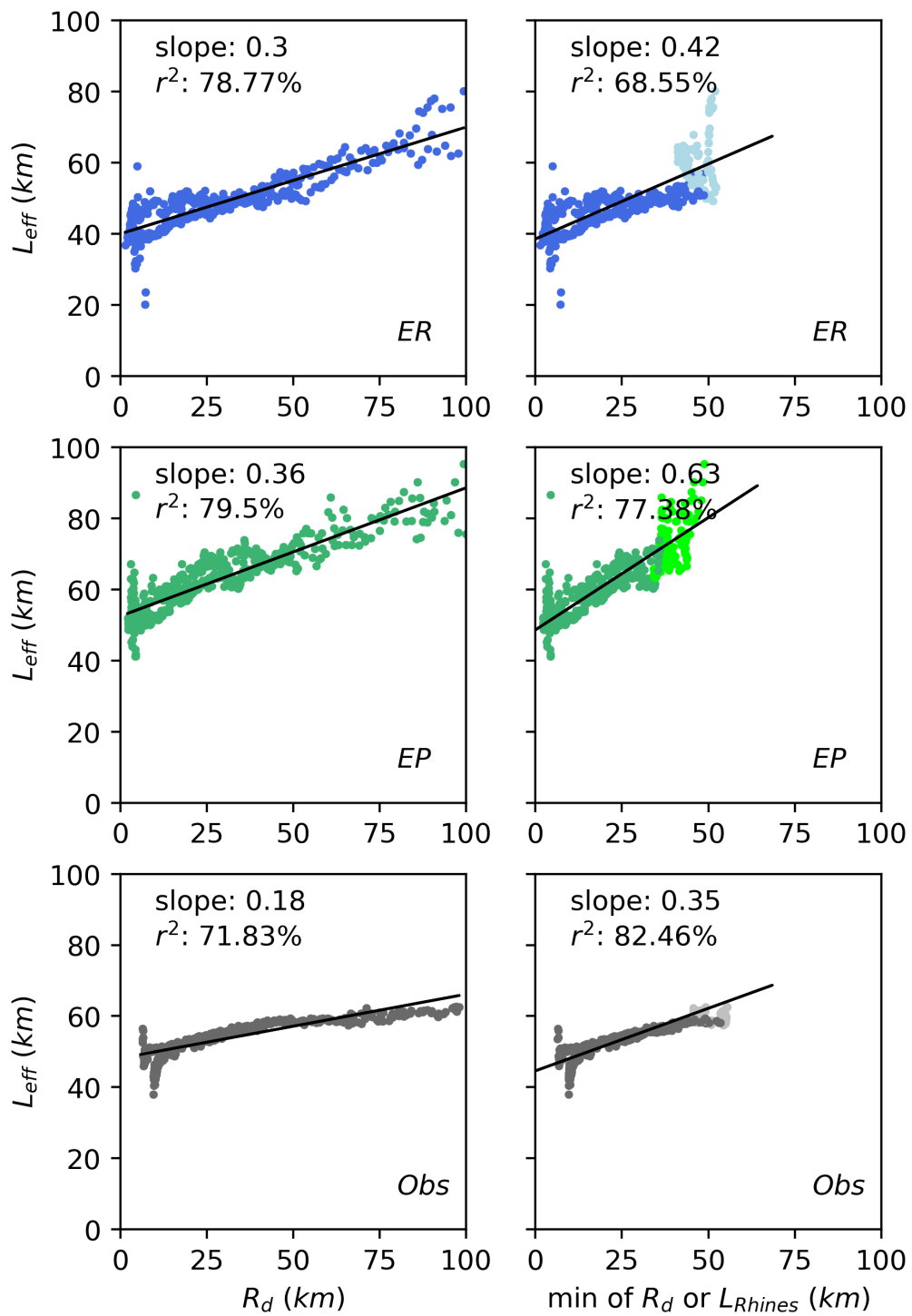


Figure A5: Same as Fig. 11 but for the effective radius L_{eff} .

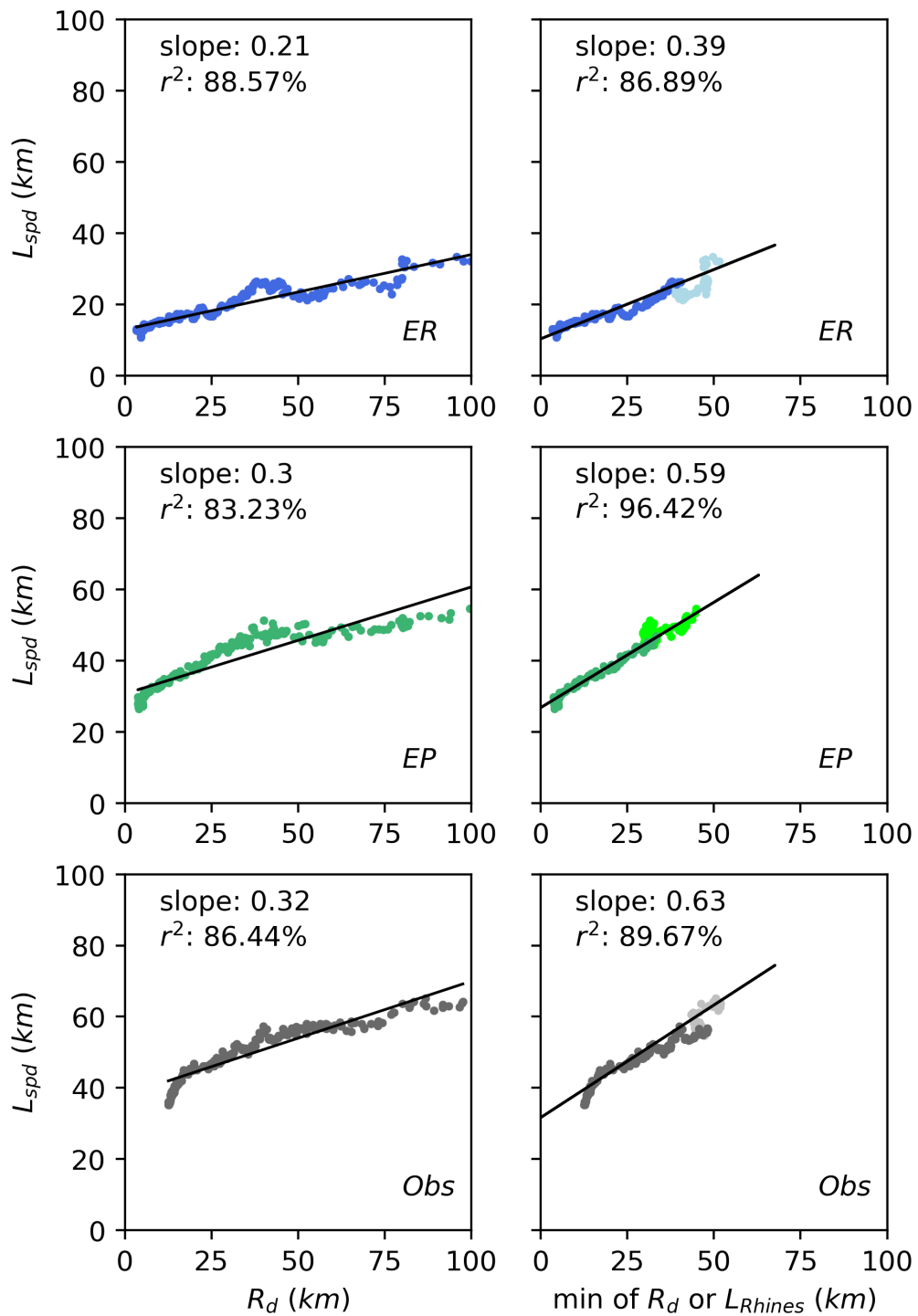


Figure A6: Same as Fig. 11 but for the North Atlantic only ($0 - 70^\circ N$, $80^\circ W - 10^\circ E$) for comparison with [22].